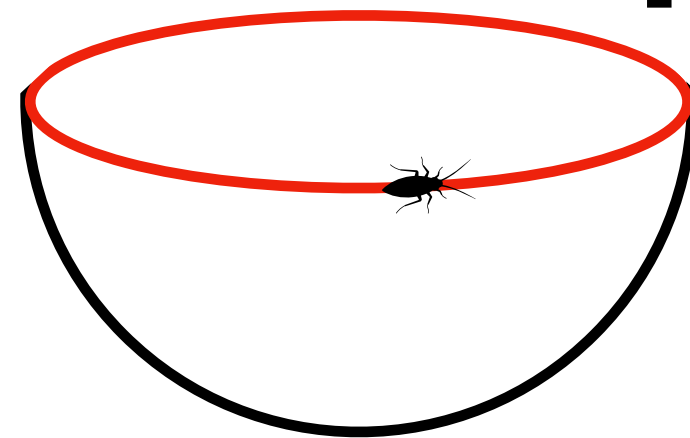
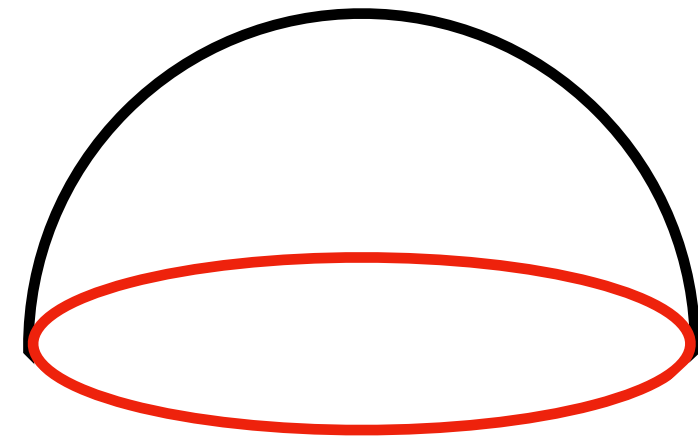


dS^4 Metamorphosis

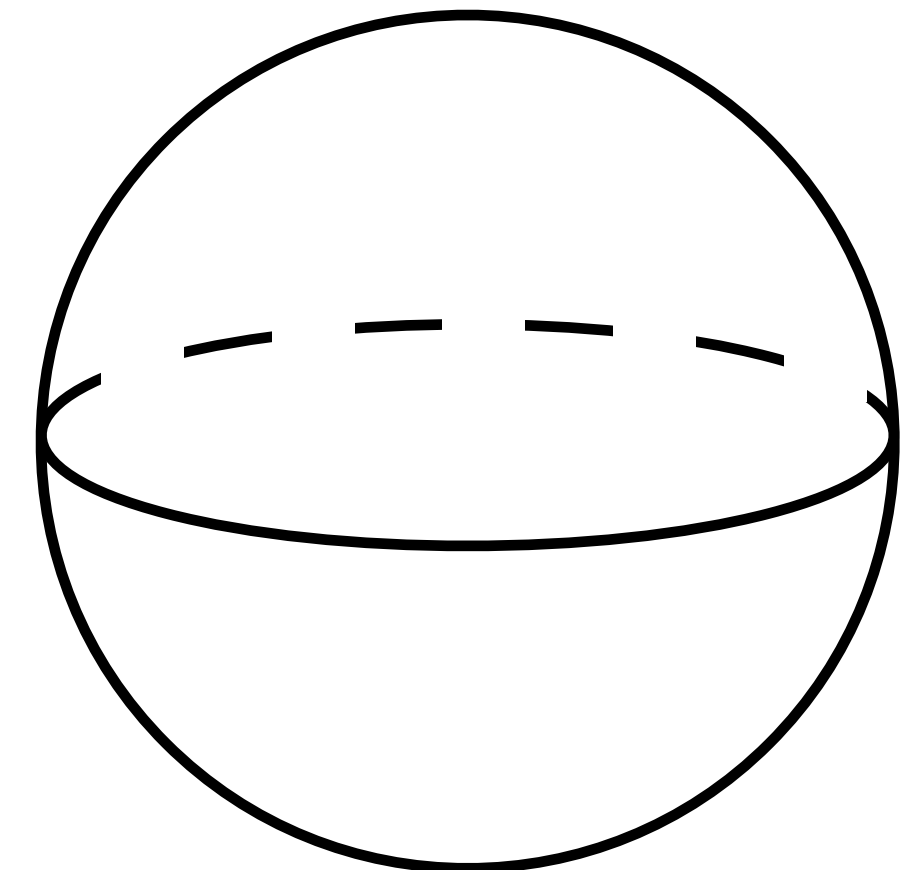


Beatrix Mühlmann

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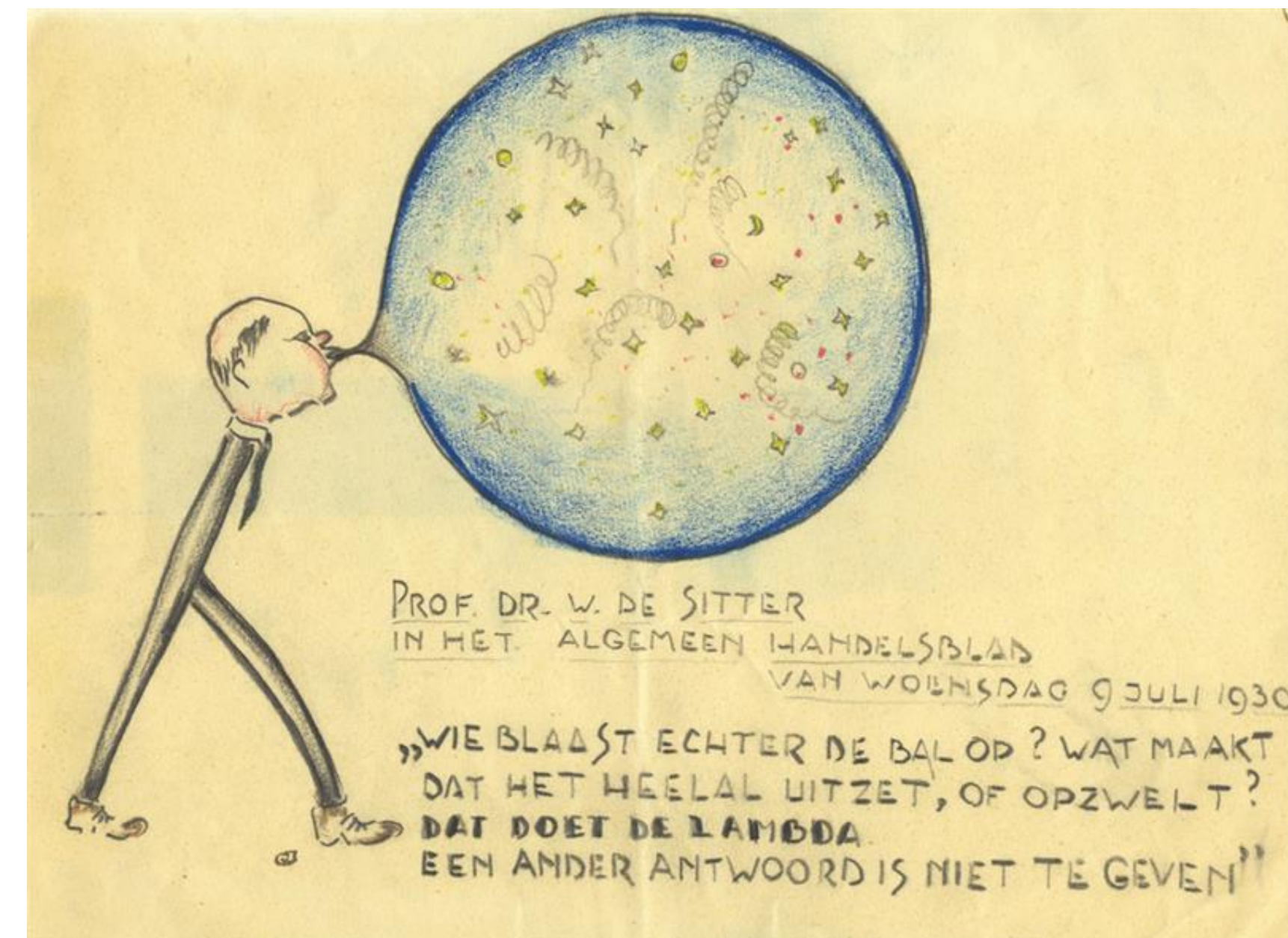
Observers, wormholes and complex saddles in cosmology

ArXiv: 2602.19812 D. Anninos, C. Baracco, BM, V. Letsios



De Sitter quantum gravity

Models of de Sitter as tractable laboratories



dS_2 Timelike Liouville theory, Sine Dilaton gravity, JT gravity,...

[...]

dS_3 Complex Liouville string as matrix model/dS frame, timelike boundaries, DSSYK,...

[...]

dS_4 Late time wavefunction in dS_4 higher spin dual to $Sp(N)$ CFT₃,...

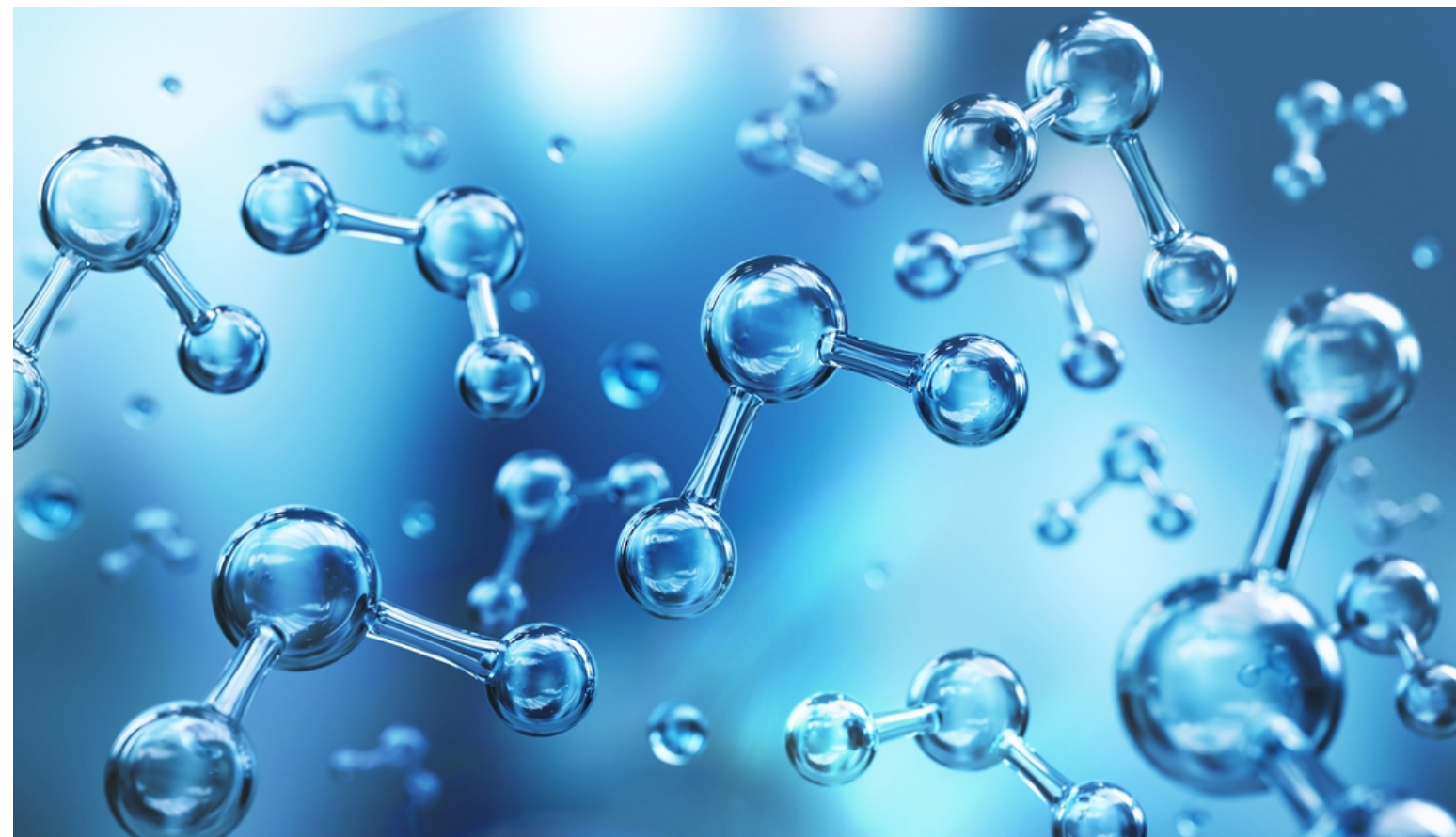
[...]

Sphere partition function

- Sphere partition function as gauge-invariant, systematically calculable “observable”
- Sphere is the Euclidean realisation of de Sitter space
- Conjecturally it encodes the de Sitter entropy

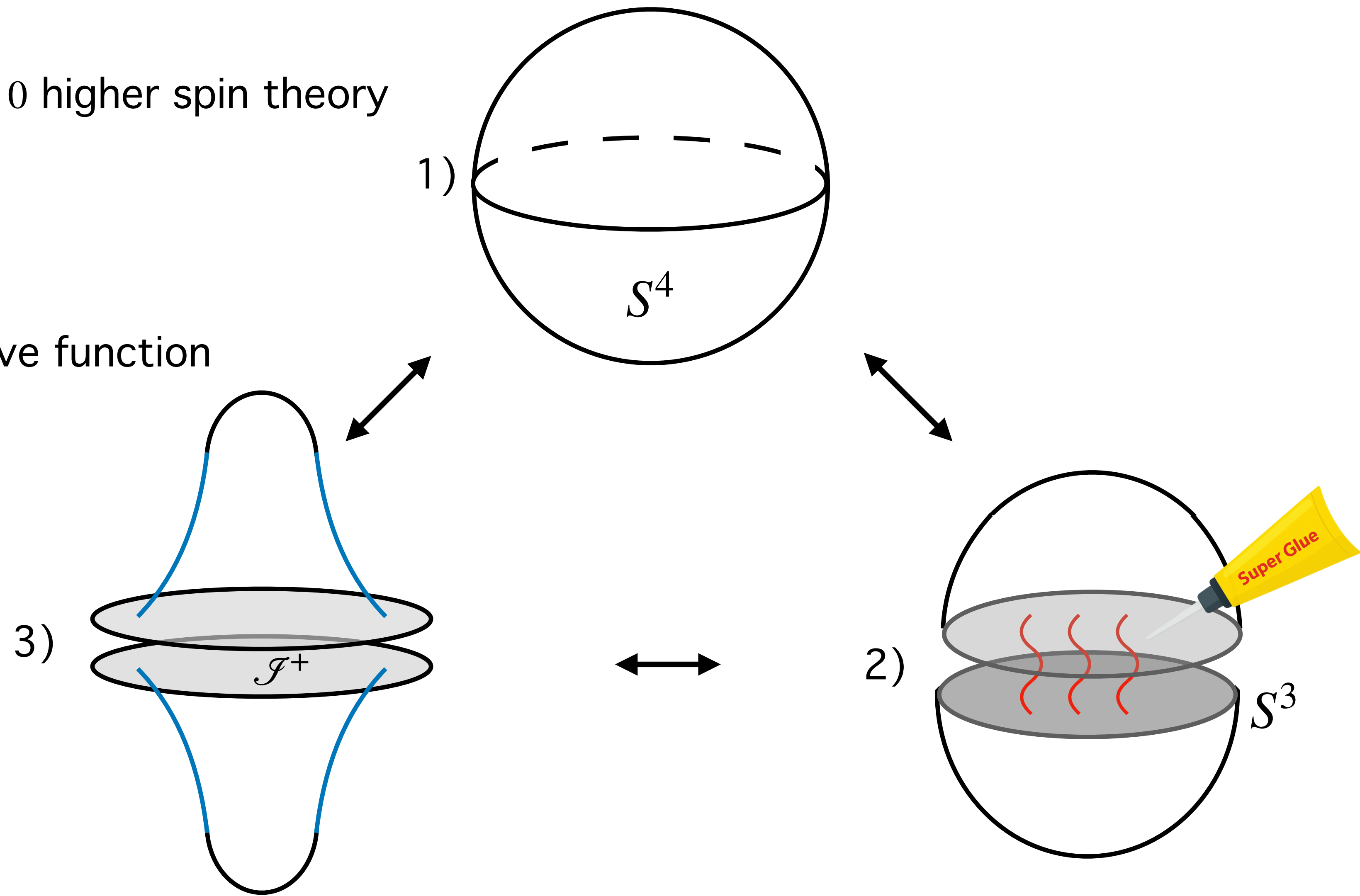
[Gibbons-Hawking,...]

What are the molecules of our Universe?



Outlook

- 0) Useful math framework
- 1) Sphere partition function of $\Lambda > 0$ higher spin theory
- 2) (Super)-gluing formula
- 3) dS/CFT and the norm of the wave function



Irreps

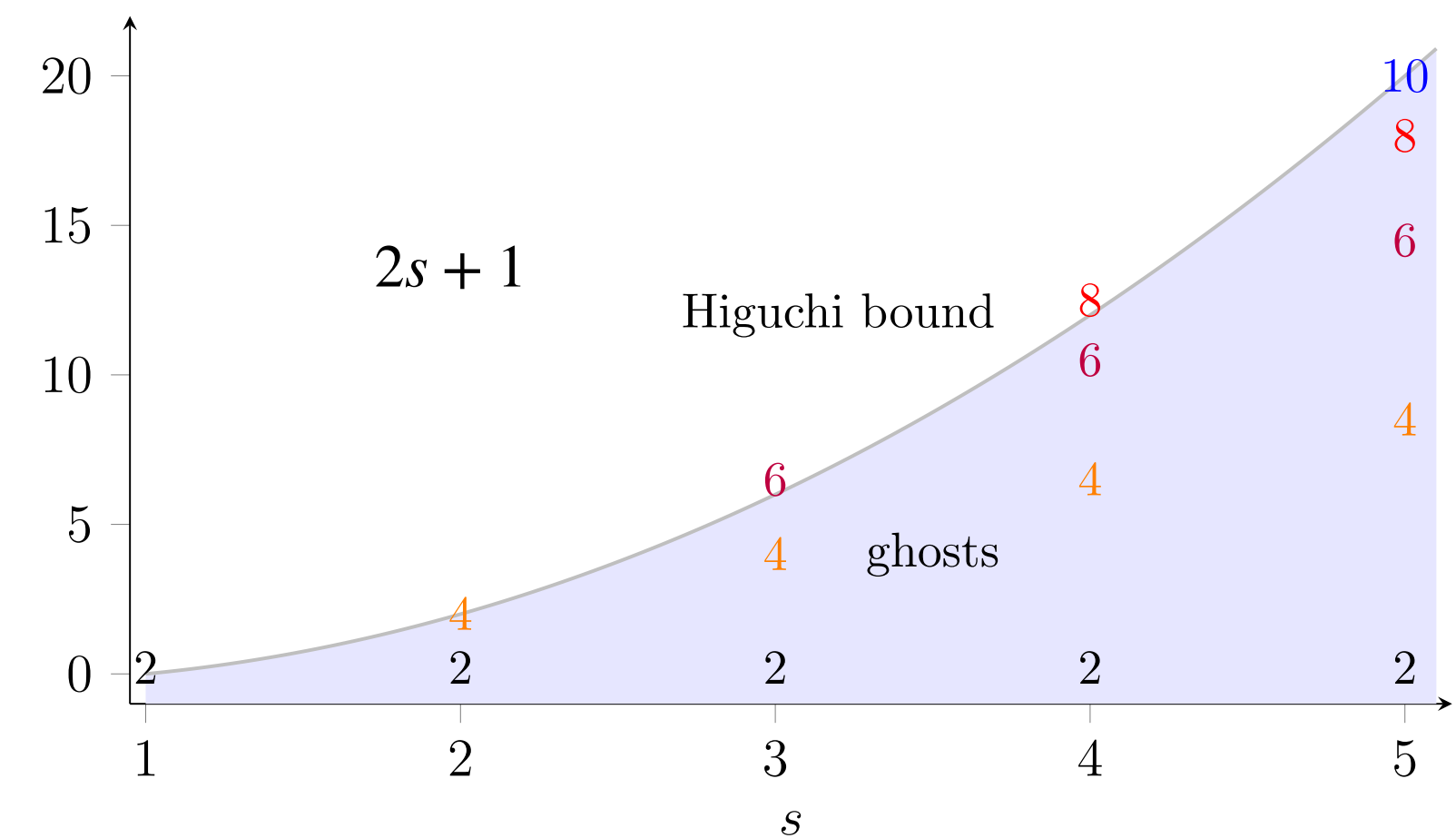
Classify free fields in dS according to the unitary irreducible representations of its isometry group

[Thomas, Newton, Dixmier, Harish-Chandra,.....]

$dS_4 \Leftrightarrow SO(1,4)$

Irrep	Range of Δ	Range of s	
π_ν	$\Delta = \frac{3}{2} + i\nu, \quad \nu \in \mathbb{R}_+$	$s = 0, 1, 2, \dots$	← Heavy fields
γ_Δ	$0 < \Delta < 3$	$s = 0$	← Light fields
	$0 < \Delta < 2$	$s \geq 1$	
$\mathcal{E}_{\Delta,0}$	$\Delta = 2 + p, \quad p \geq 1$	$s = 0$	
$\mathcal{D}_{t,s}^\pm$	$\Delta = 2 + t, \quad t = 0, 1, \dots, s-1$	$s \geq 1$	← Gauge fields $m^2 \ell^2$

Partially massless fields



Characters

Harish-Chandra character for $SO(1, d + 1)$

[Harish-Chandra, Hirai,...]

$$\chi_{\Delta, s}(\mathbf{t}) = \text{tr}_{\Delta, s} e^{-itH} = \int_{\mathbb{R}} d\omega \rho_{\Delta, s}(\omega) e^{-i\omega t}, \quad H \in \mathfrak{so}(1, 1)$$

Example 1 Conformally coupled scalar in dS_4 ($m^2 = 2 \Leftrightarrow m^2 = \Delta(3 - \Delta)$)

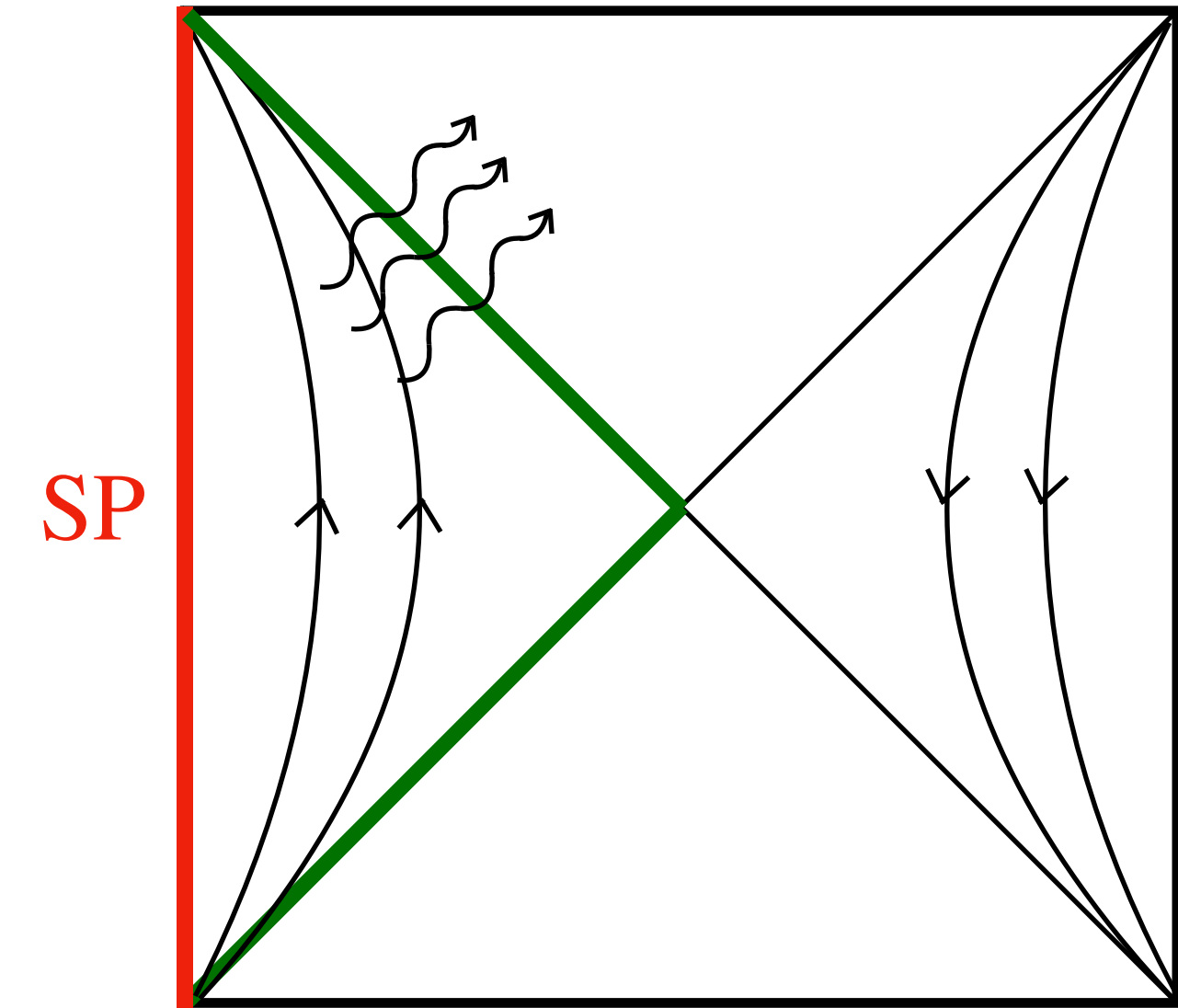
$$\chi_{1,0}(t) = \frac{e^{-2t} + e^{-t}}{(1 - e^{-t})^3} = e^{-t} + 4e^{-2t} + 9e^{-3t} + 16e^{-4t} + \dots$$

Degeneracies of quasinormal modes

[Lopez-Ortega,..., Anninos-Denef-Law-Sun,...]

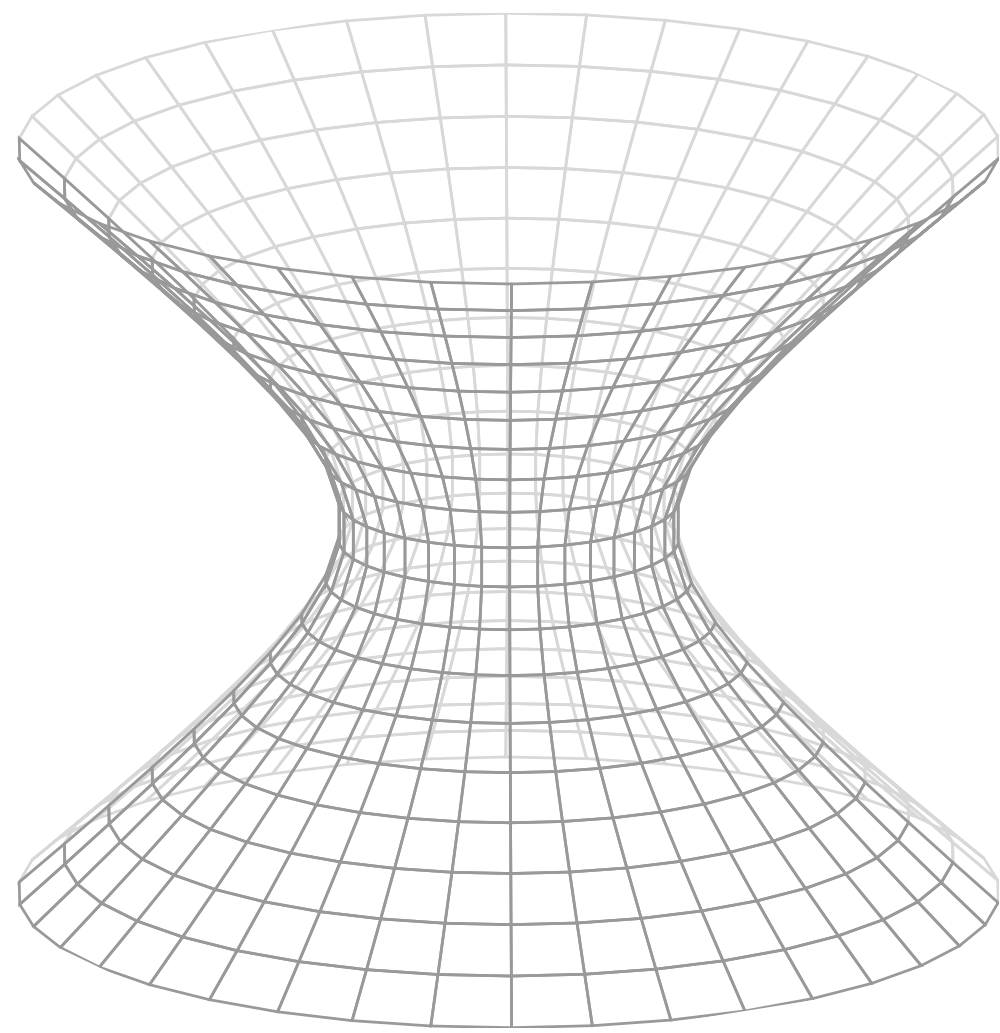
Example 2 Graviton in dS_4

$$\chi_{\text{grav}}(\mathbf{t}) = \chi_{\text{bulk}}(\mathbf{t}) - \chi_{\text{edge}}(\mathbf{t}) \quad \chi_{\text{bulk}}(\mathbf{t}) = 2 \frac{5e^{-3t} - 3e^{-4t}}{(1 - e^{-t})^3}, \quad \chi_{\text{edge}}(\mathbf{t}) = \frac{10e^{-2t} - 2e^{-3t}}{(1 - e^{-t})}$$



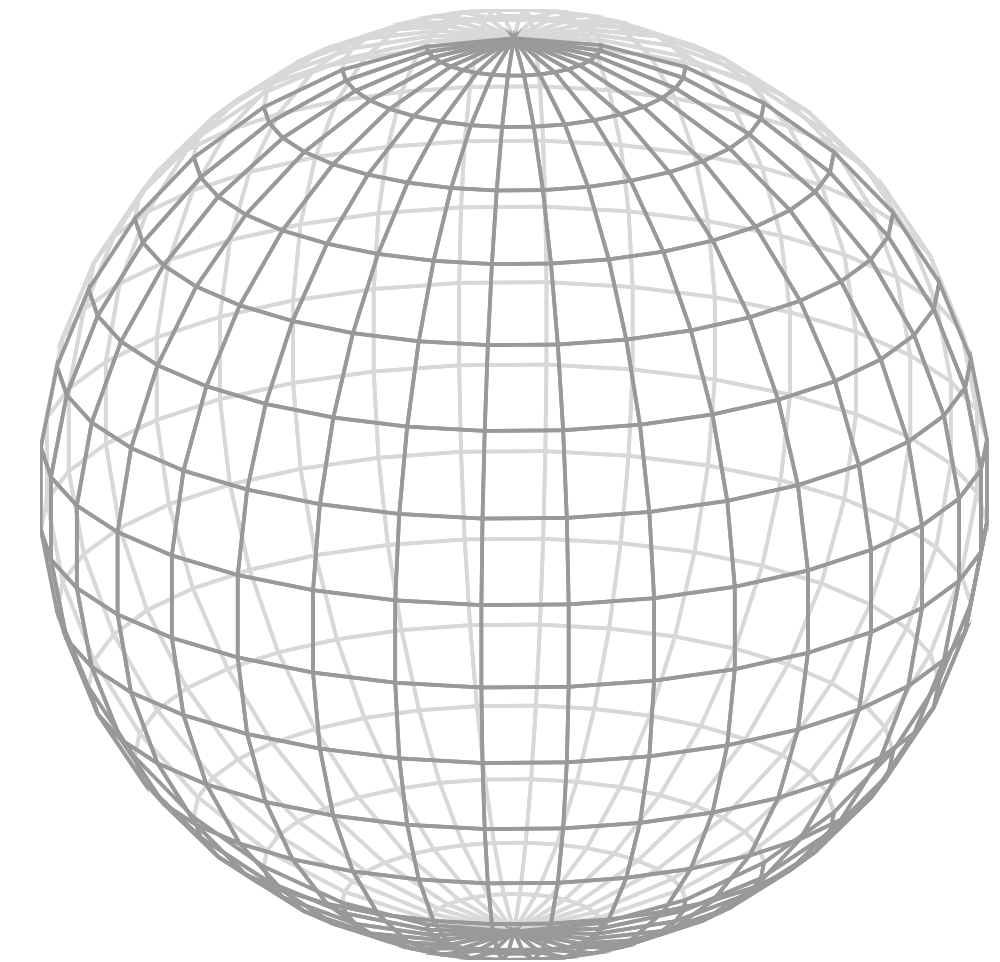
1-loop

Lorentzian Harish-Chandra characters



[Anninos-Denef-Law-Sun,...]

Euclidean one-loop sphere partition functions



$$\text{Example: } \log \mathcal{Z}_{\text{P.I.}}^{(1)} = \log \det^{-1/2} \left(-\nabla_{S^4}^2 + 2 \right) = \int_0^\infty \frac{dt}{2t} \frac{1 + e^{-t}}{1 - e^{-t}} \chi_{1,0}(t)$$

Minimal higher spin theory

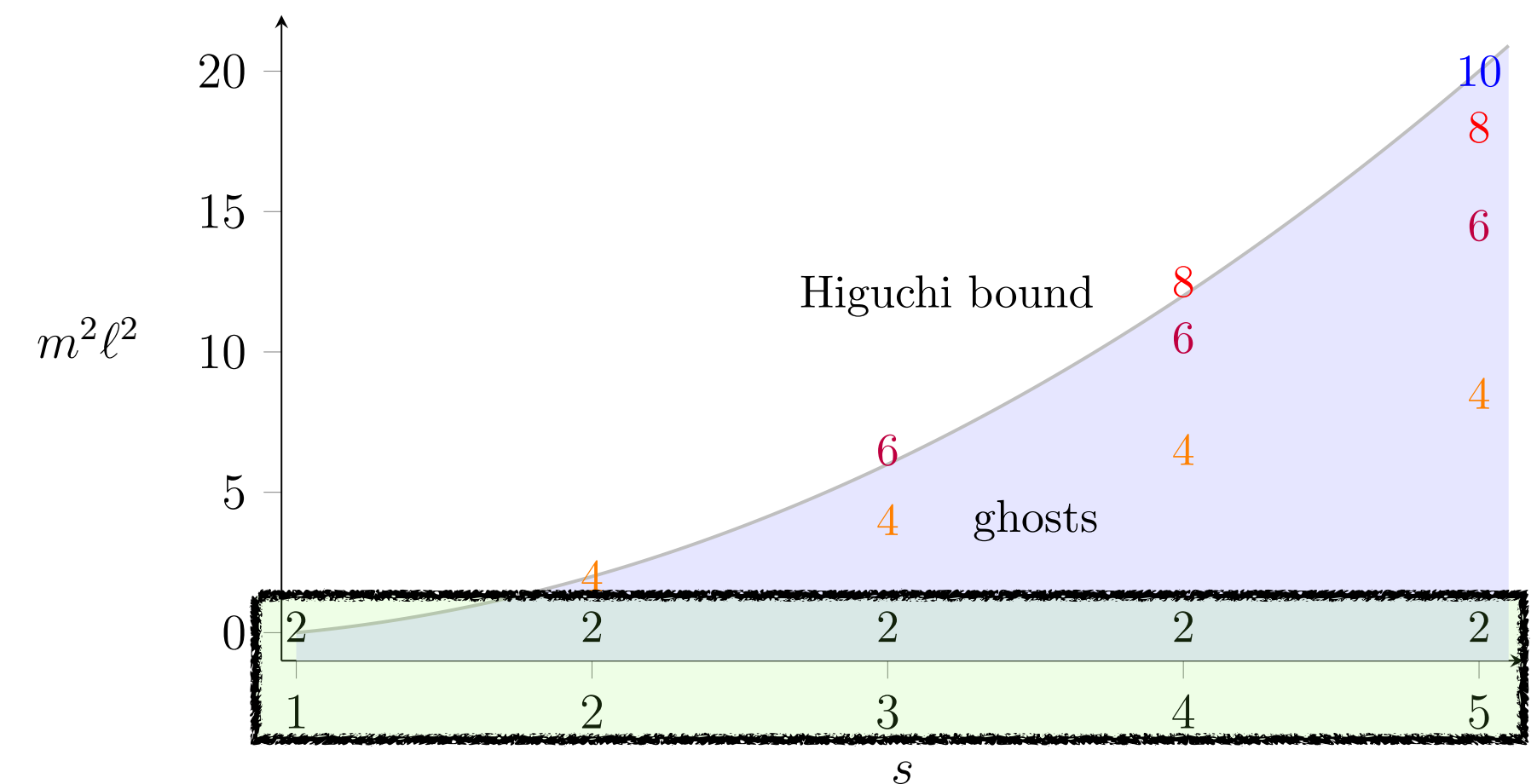
Minimal higher spin theory consists of $\{b, b_{\nu_1\nu_2}, b_{\nu_1\nu_2\nu_3\nu_4}, \dots\}$
 [Fronsdal, Fradkin-Vasiliev, Vasiliev,...]

Conformally coupled scalar

Highest depth partially massless fields

The Fronsdal fields $b_{\nu_1\dots\nu_s}$

- totally symmetric & $b^{\nu\mu}_{\nu\mu\nu_5\dots\nu_s} = 0$
- are invariant under $b_{\nu_1\dots\nu_s} \rightarrow b_{\nu_1\dots\nu_s} + \nabla_{(\mu} \zeta_{\nu_2\dots\nu_s)}$, $\zeta^{\mu}_{\mu\nu_3\dots\nu_{s-3}} = 0$
- have two degrees of freedom
- are highest depth PMF with spin s and $\Delta = s + 1$



Higher spin at 1-loop

What is the microscopic completion of $\Lambda > 0$ higher spin on S^4 ?

$$\log Z_{\text{hs}}^{(1)} = \int_0^\infty \frac{dt}{2t} \frac{1+e^{-t}}{1-e^{-t}} \sum_{s \in 2\mathbb{N}} \chi_{1+s,s}(t) = \int_0^\infty \frac{dt}{2t} \frac{1+e^{-t}}{1-e^{-t}} \left(\frac{-e^{-t}}{(1-e^{-t})^2} + 2 \frac{e^{-\frac{3}{2}t} + e^{-\frac{t}{2}}}{(1-e^{-t})^2} \right)$$
$$= \frac{\zeta(3)}{8\pi^2} + 2 \left(\frac{3\zeta(3)}{16\pi^2} - \frac{1}{8} \log 2 \right)$$

 The LHS organises itself into 3D characters

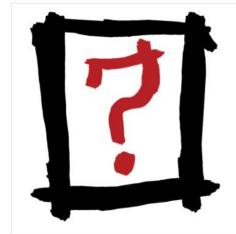
 Conformal higher spin theory

 Conformally coupled scalar in dS_3

Sphere partition function

The sphere partition function also needs the on-shell action which famously is unknown

$$Z_{\text{hs}}[S^4] \approx e^{-S_{\text{EdS}}^{(N)}} Z_{\text{h.s.}}^{(1)}$$



❗ AdS/CFT: Predicts $S_{\text{EAdS}}^{(N)}$ with $N \approx (\Lambda G_N)^{-1}$

❗ Conjecture: $-S_{\text{EdS}}^{(N)} = -2S_{\text{EAdS}}^{(-N)} > 0$

Together this leads to
$$Z_{\text{hs}}[S^4] \approx \frac{(-i)^{\mathcal{P}}}{\text{vol}_N(G_{\text{hs}})} e^{2N \left(\frac{1}{8} \log 2 - \frac{3\zeta(3)}{16\pi^2} \right) + \frac{\zeta(3)}{8\pi^2}}$$

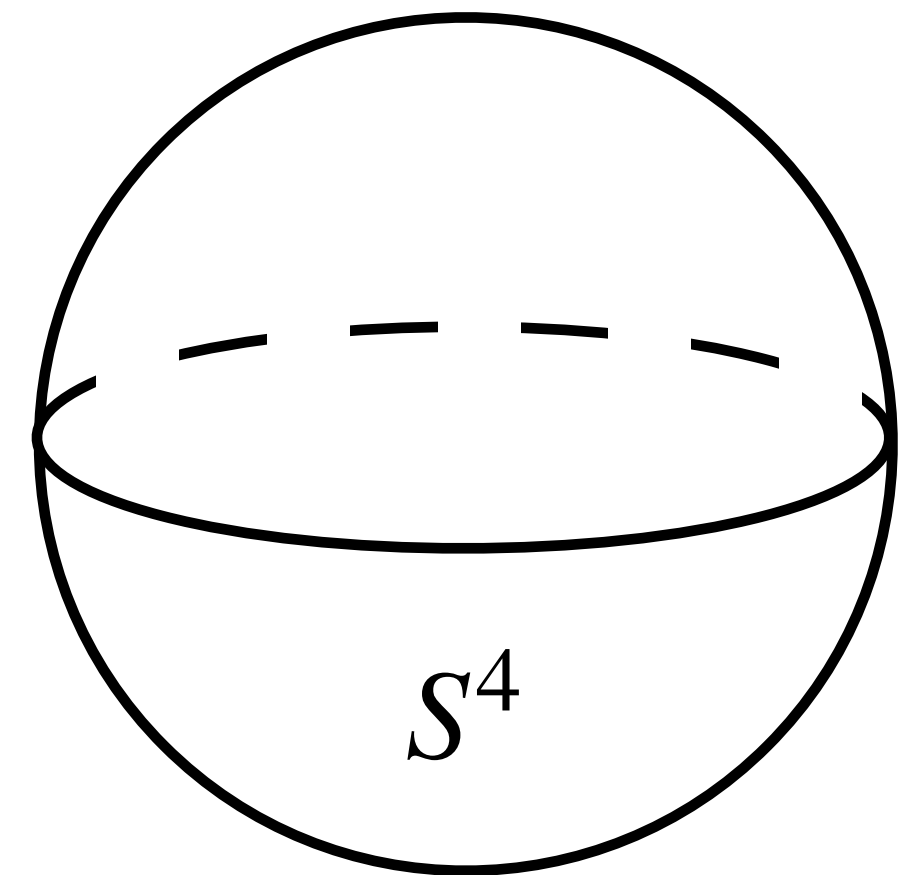
\mathcal{P}

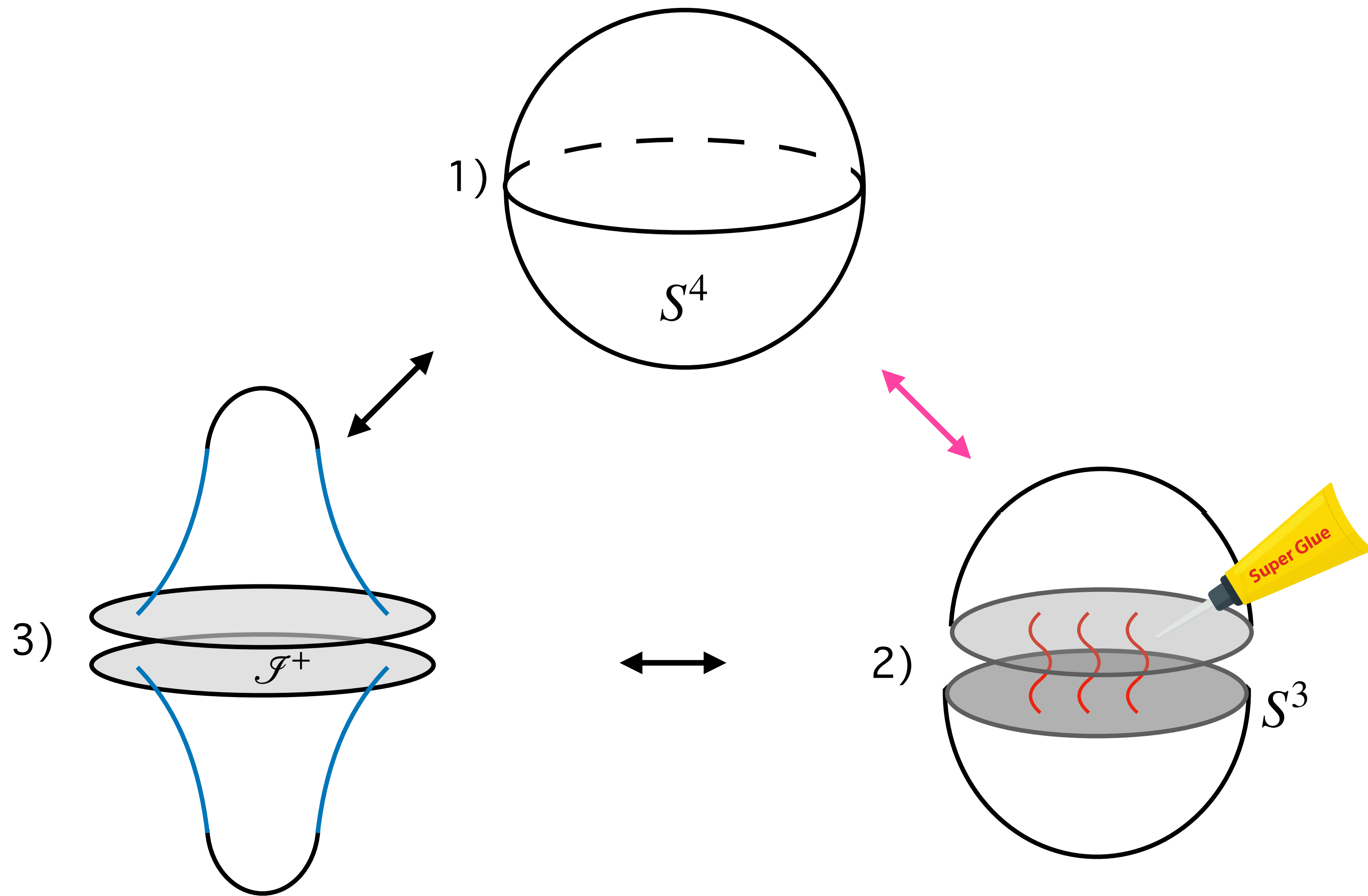
The phase cancels in some cases

[Giombi-Sun, Anninos-Baracco-BM-Letsios,...]

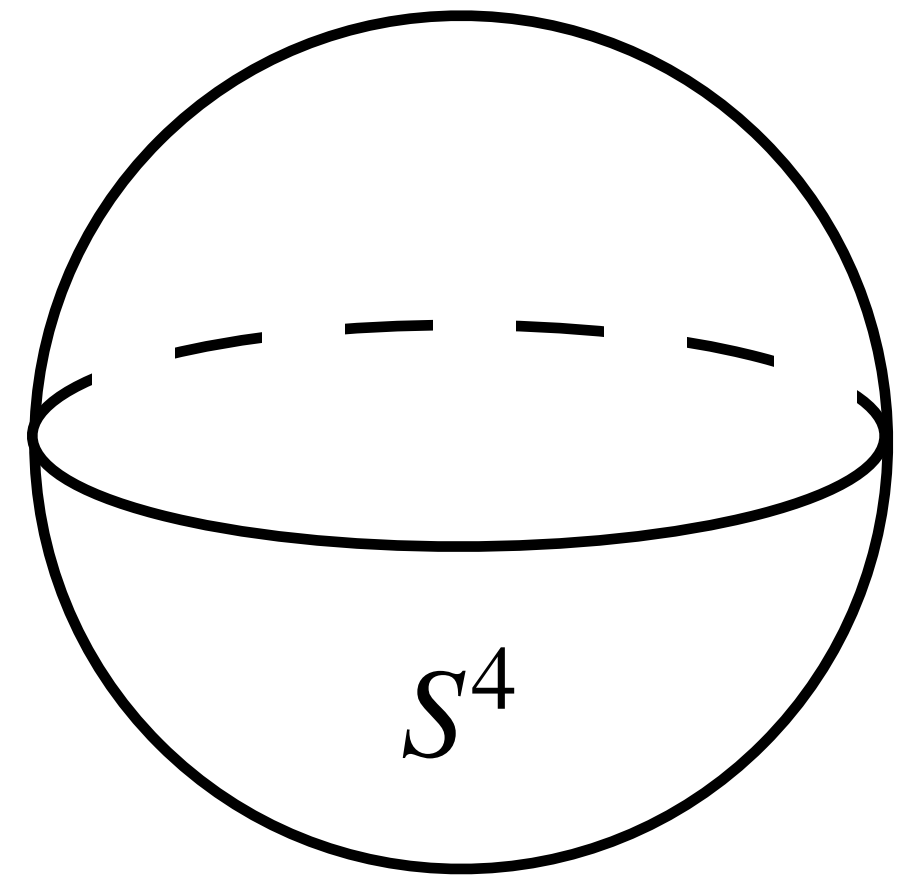
G_{hs}

Generated by spin s Killing tensors

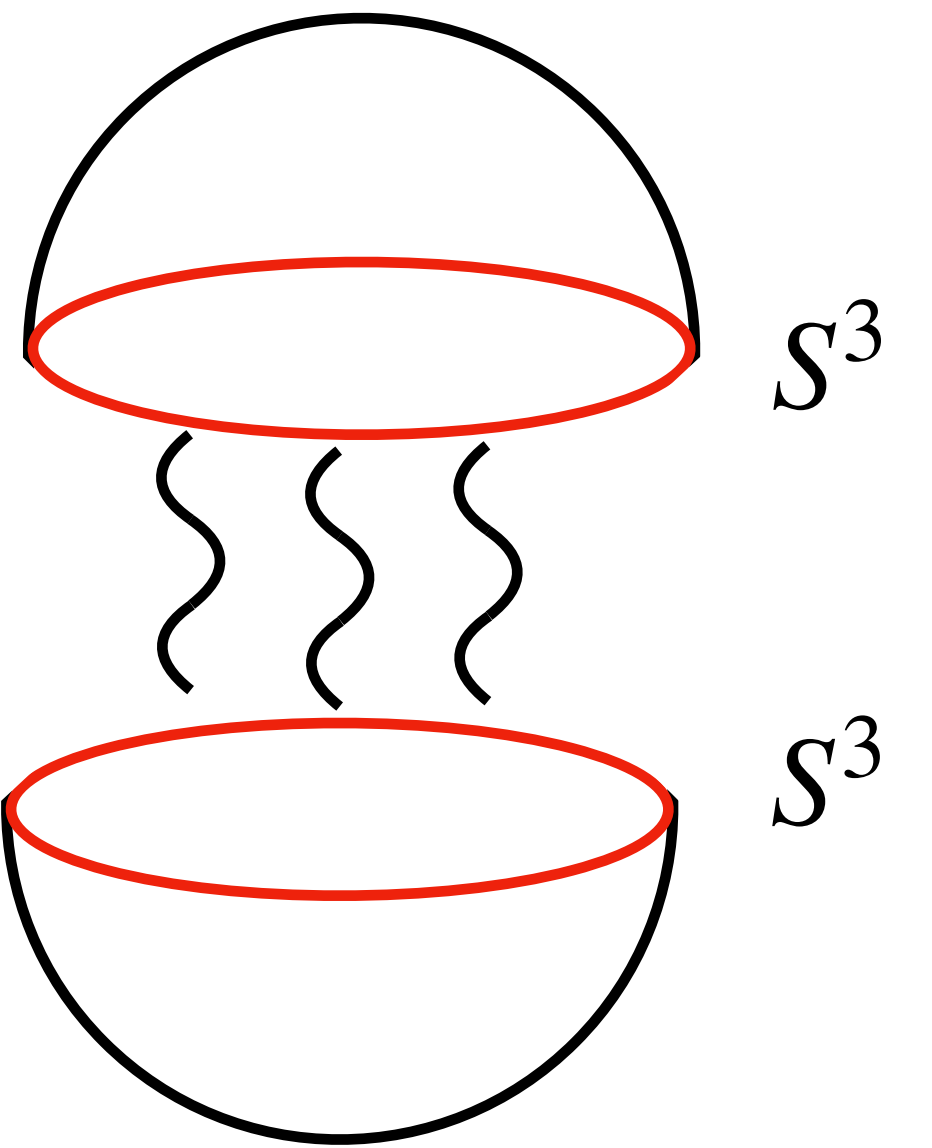




Conjecture



A diagram of a sphere labeled S^4 . The sphere is drawn with a solid line for the front half and a dashed line for the back half to indicate its three-dimensional nature.

$$= Z_{\text{hs}}[S^4] = \frac{(-i)^{\mathcal{P}}}{\text{vol}_N(G_{\text{hs}})} \int [\mathcal{D}B] \left| Z_{Sp(N)}[B] \right|^2 =$$


A diagram showing two spheres, each labeled S^3 , one above the other. The top and bottom boundaries of both spheres are highlighted with red lines. Three wavy lines connect the top and bottom boundaries, representing interactions or fields between the two spheres.

$Z_{Sp(N)}$ Partition function of $Sp(N)$ vector model build from *anticommuting* scalars

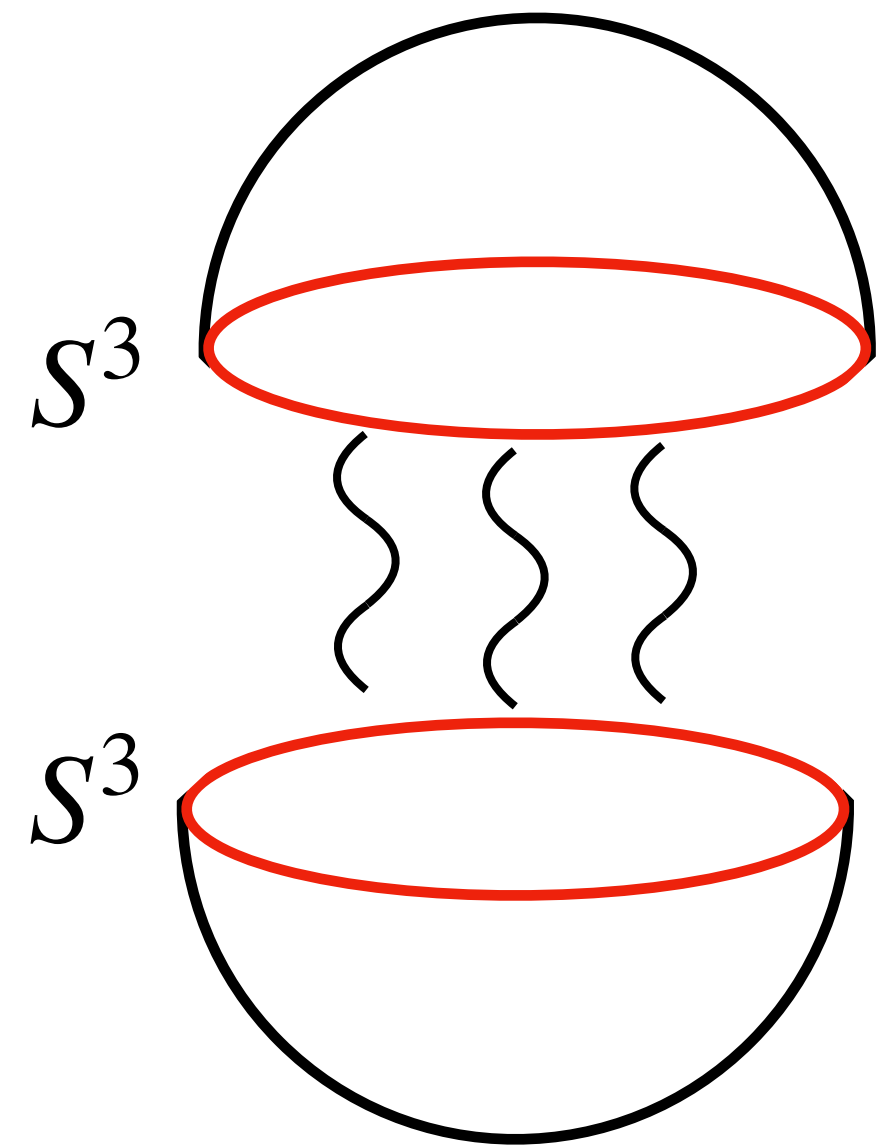
$$Z_{Sp(N)}[B] = \int [\mathcal{D}\chi_I] e^{-\int_{S^3} \chi^I (-\nabla_c^2) \chi^J \Omega_{IJ} + b \mathcal{O} + b^{ij} \mathcal{O}_{ij} + \dots}$$

$$\begin{aligned} \mathcal{O} &= \chi^I \chi^J \Omega_{IJ} & I = 1, \dots, N \\ \mathcal{O}_{ij} &= \chi^I \partial_i \partial_j \chi^J \Omega_{IJ} \\ &\vdots \end{aligned}$$

B bilocal field which sources the $Sp(N)$ invariant operators

G_{hs} 4D Vasiliev algebra = 3D Eastwood algebra
 ($SO(1,4) = dS_4$ isometry group and 3D Euclidean conformal group)

Gluing formula

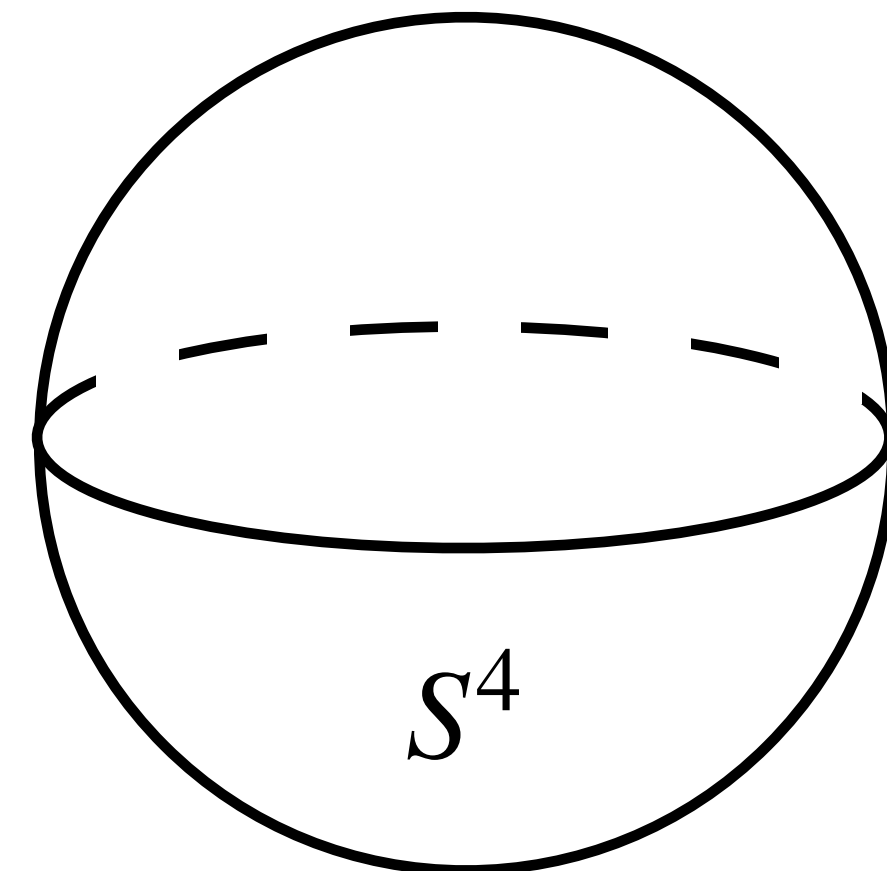


$$= \frac{(-i)^{\mathcal{P}}}{\text{vol}_N(G_{\text{hs}})} \int [\mathcal{D}B] \left| Z_{Sp(N)}[B] \right|^2 = \frac{(-i)^{\mathcal{P}}}{\text{vol}_N(G_{\text{hs}})} \int [\mathcal{D}B] \det^N(-\nabla_c^2 + B)$$

$$\approx \frac{(-i)^{\mathcal{P}}}{\text{vol}_N(G_{\text{hs}})} \times \boxed{\det^N(-\nabla_c^2)} \times \boxed{\text{higher spin two-point function}}$$

[Giombi-Klebanov-Pufu-Safdi-Tarnopolsky,...]

$$= \frac{(-i)^{\mathcal{P}}}{\text{vol}_N(G_{\text{hs}})} e^{2N(\frac{1}{8} \log 2 - \frac{3\zeta(3)}{16\pi^2}) + \frac{\zeta(3)}{8\pi^2}} =$$



Extension

Integer + half-integer higher spin fields

$$\{b, b_{\nu_1}, b_{\nu_1\nu_2}, b_{\nu_1\nu_2\nu_3}, \dots\}$$

Conformally coupled scalar Highest depth PMF $s = 1, 2, \dots$

+

$$\{\psi_\alpha, \psi_{\alpha\mu_1}, \psi_{\alpha\mu_1\mu_2}, \dots\}$$

Massless Dirac Complex massless gauge fields $s = \frac{3}{2}, \frac{5}{2}, \dots$

Irrep	Range of Δ	Range of s
π_ν	$\Delta = \frac{3}{2} + i\nu, \nu \in \mathbb{R}_+$	$s = 0, 1, 2, \dots$
γ_Δ	$0 < \Delta < 3$	$s = 0$
	$0 < \Delta < 2$	$s \geq 1$
$\mathcal{D}_{t,s}^\pm$	$\Delta = 2 + t, t = 0, 1, \dots, s-1$	$s \geq 1$

Irrep	Range of Δ	Range of s
π_ν	$\Delta = \frac{3}{2} + i\nu, \nu \in \mathbb{R}_+$	$s = \frac{1}{2}, \frac{3}{2}, \frac{5}{2}, \dots$
	X	
$\mathcal{D}_{t,s}^\pm$	$\Delta = 2 + t, t = -\frac{1}{2}, \frac{1}{2}, \dots, s-1$	$s = \frac{1}{2}, \frac{3}{2}, \frac{5}{2}, \dots$

Superglubits

Combining the on-shell action and the 1-loop contribution leads to

$$Z_{\mathcal{N}=2 \text{ hs}}^{(N)}[S^4] \approx \frac{1}{\text{vol}_N(G_{\text{s.hs}})} \times e^{-S_{\text{EdS}}^{(N)}} \times Z_{\mathcal{N}=2 \text{ hs}}^{(1)} = 2^N \times \frac{N^{-\frac{1}{16}}}{\text{vol}(G_{\text{s.hs}})}$$

▣ $\text{vol}_N(G_{\text{s.hs}}) = N^{\frac{1}{16}} \text{vol}(G_{\text{s.hs}})$

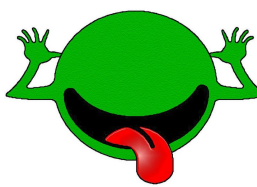
▣ $-S_{\text{EdS}}^{(N)} = N \log(2) - \frac{\zeta(3)}{4\pi^2}$

▣ $\log Z^{(1)} = \frac{\zeta(3)}{4\pi^2}$

↑
Σ Integer+half-integer KT/2

Ratio is independent of $G_{\text{s.hs}}$

$$r_{N,M} \equiv \frac{Z_{\mathcal{N}=2 \text{ hs}}^{(N)}[S^4]}{Z_{\mathcal{N}=2 \text{ hs}}^{(M)}[S^4]}$$

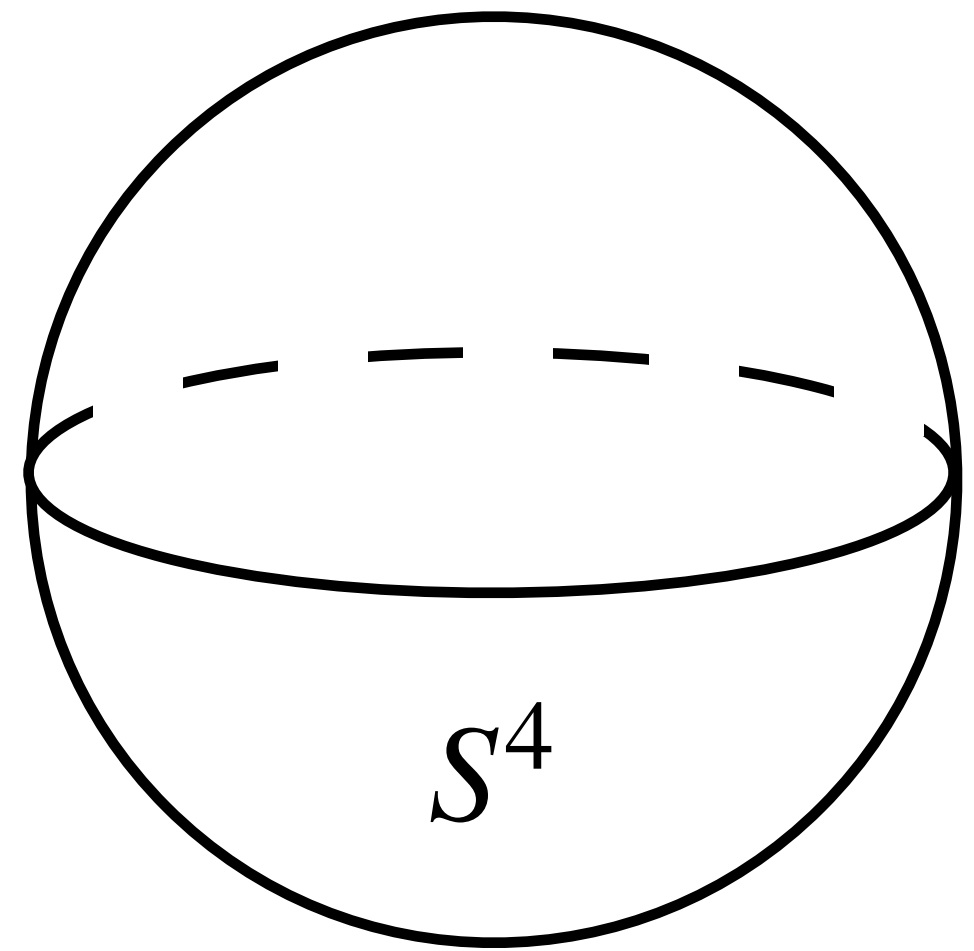
$N = (432) , \quad M = (432)2^{16}$
 $\Rightarrow \log(r_{N,M}) = 10^{42}$ 

Supergluing formula

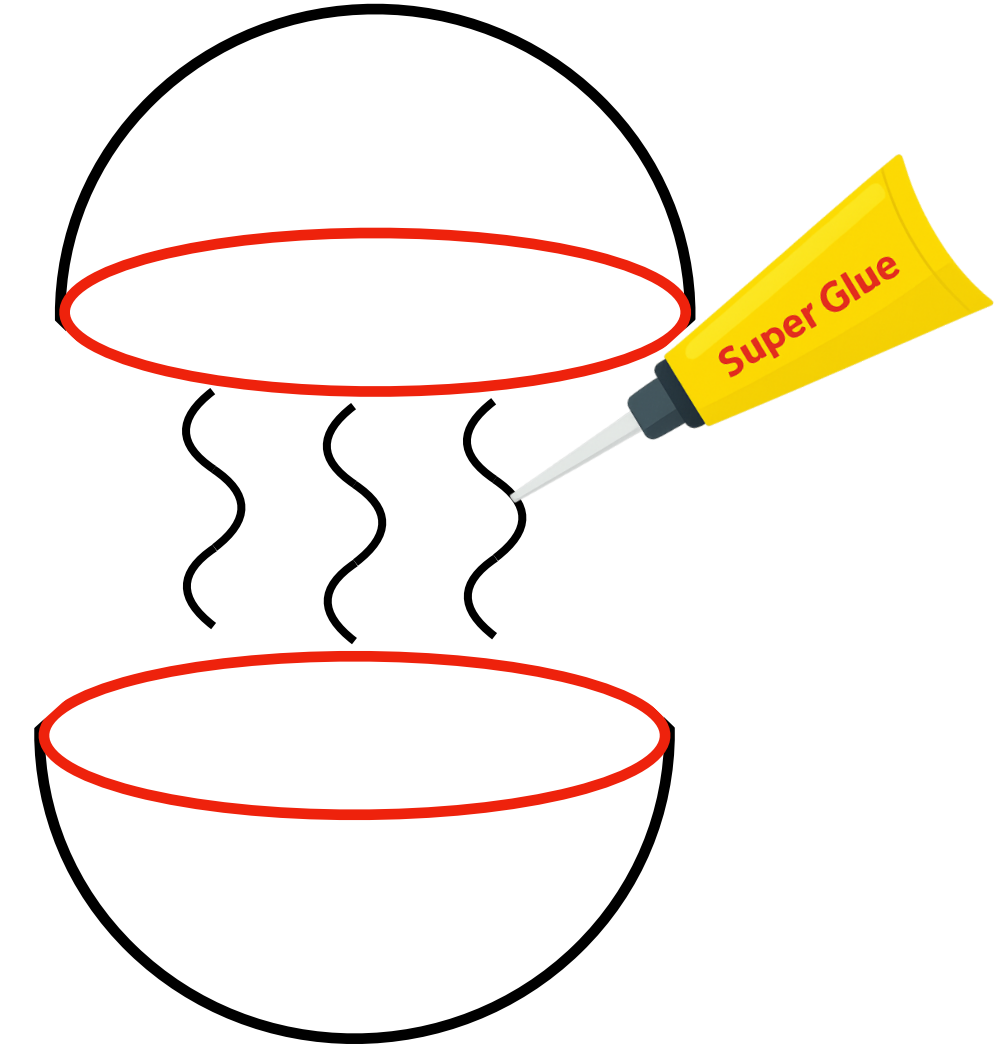
$$Z_{\mathcal{N}=2 \text{ hs}}[S^4] = \frac{1}{\text{vol}_N(G_{\text{s.hs}})} \int [\mathcal{D}B_s] \left| Z_{U(N)}[B_s] \right|^2$$

$Z_{U(N)}$ Partition function of N free *anticommuting* complex scalars & N free *commuting* Dirac spinors

$$Z_{U(N)}[B_s] = \int \underset{\substack{\uparrow \\ \text{chiral superfield}}}{[\mathcal{D}\Phi_I]} \underset{\substack{\uparrow \\ \text{anti-chiral superfield}}}{[\mathcal{D}\tilde{\Phi}_I]} e^{-\int dz \tilde{\Phi}_I(\tilde{z})\Phi_I(z) + \int dz_1 dz_2 \tilde{\Phi}_I(\tilde{z}_1) \underset{\substack{\uparrow \\ \text{bi-local superfield}}}{B_s(z_1; \tilde{z}_2)} \Phi_I(z_2)}$$



=



Outlook

■ Gluing formula for general dimension? [WIP Anninos-Baracco-BM-Letsios...]

■ $2^N = \lim_{\beta \rightarrow 0} (e^{\beta\omega/2} + e^{-\beta\omega/2})^N \Rightarrow$ microscopic picture?

■ Beyond one-loop

■ Timelike Toda as a 2D toy model [WIP BM...]

■ $\mathcal{N} = 2$ Liouville [WIP BM- Zikopolous...]

⋮

Relation to dS/CFT

Higher spin in Lorentzian dS₄

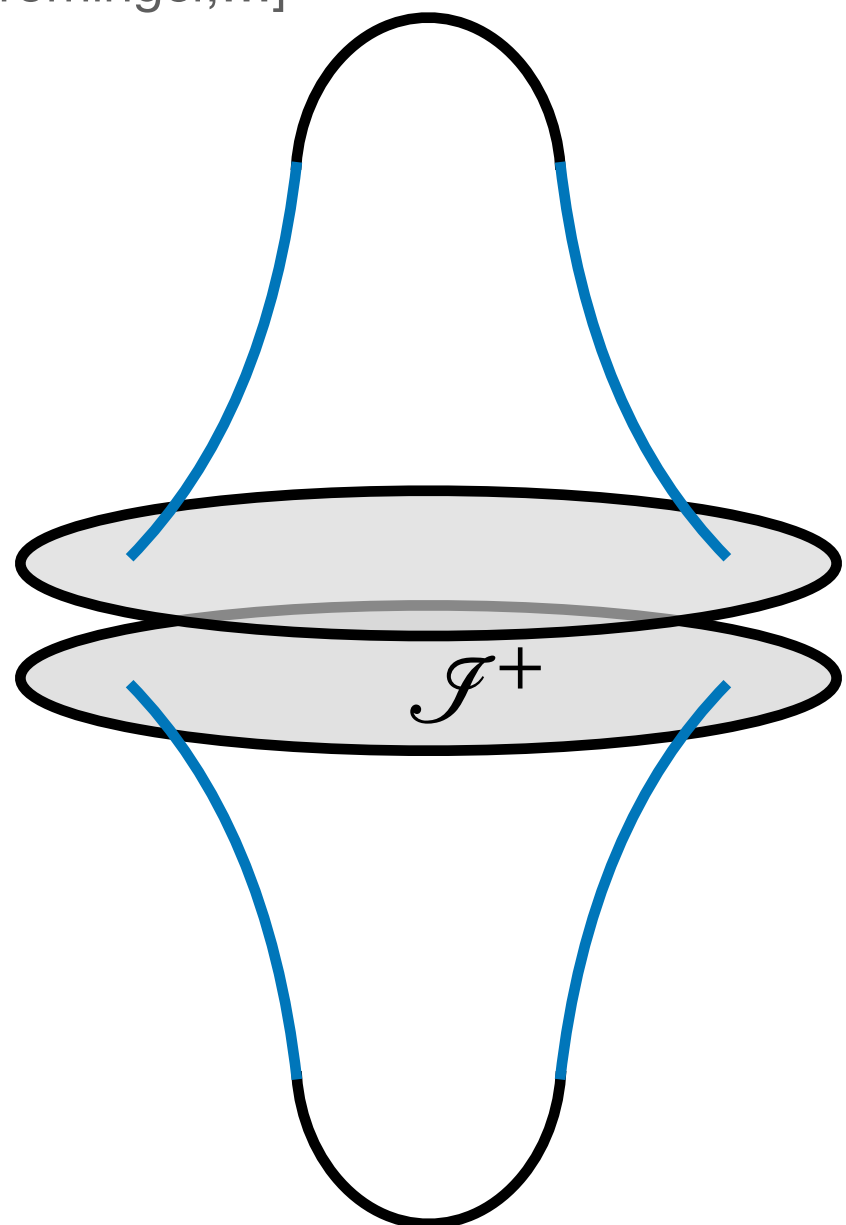
A different perspective:

dS/CFT: [Maldacena,...] HH wave function ↔ CFT partition function at \mathcal{I}^+

dS isometries become conformal transformations at \mathcal{I}^+

Higher spin dS/CFT: [Anninos-Hartman-Strominger,...]

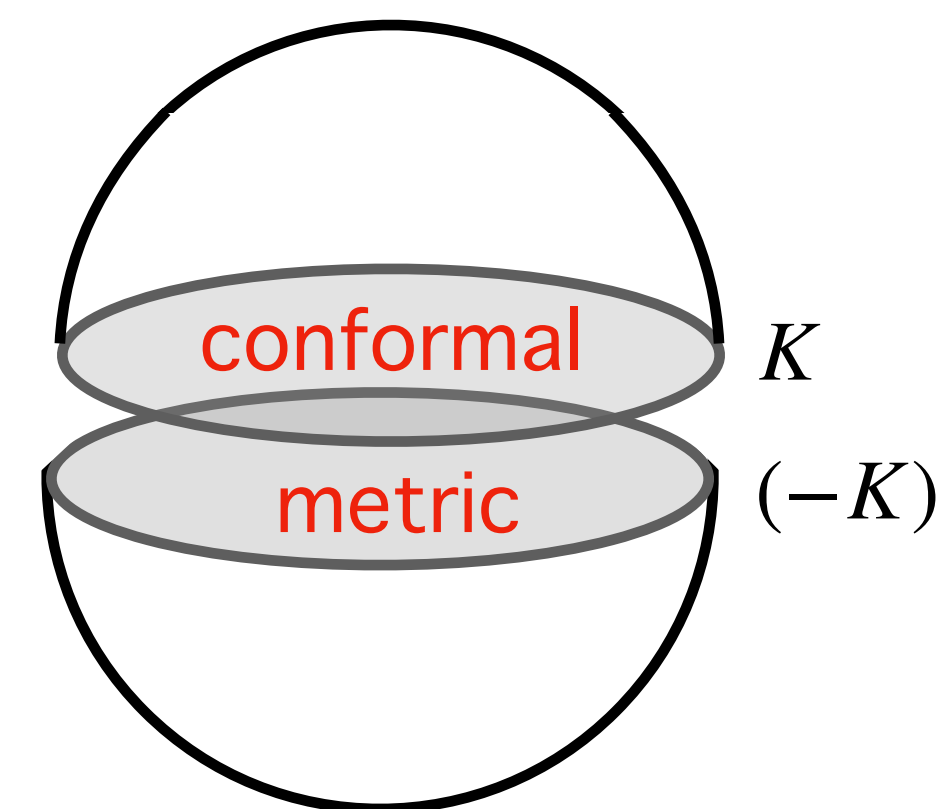
$$\Psi_{\text{HH}}[B_\mu] = Z_{Sp(N)}[B_i]$$



$$\langle \Psi_{\text{HH}} | \Psi_{\text{HH}} \rangle$$

$\stackrel{?}{=}$

$$Z_{\text{hs}}[S^4]$$

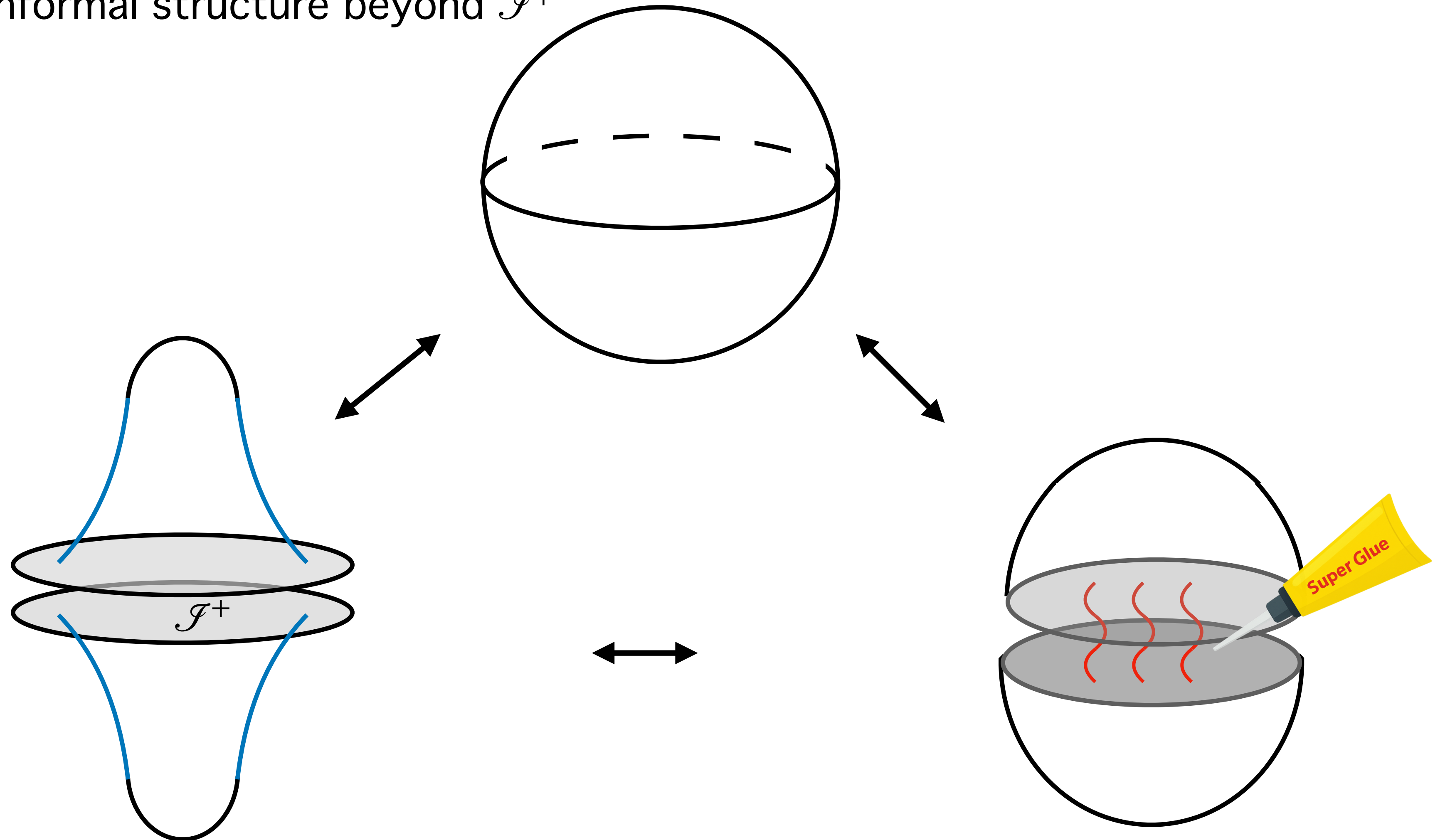


Norm at finite boundary with York time K as time parameter

[WIP Blommaert-BM-Verlinde...]

Summary

- 1) We developed a machinery to access the $\Lambda > 0$ higher spin sphere partition function
- 2) Gluing formula reveals conformal structure beyond \mathcal{I}^+
- 3) Hints of a norm?



Thank you!