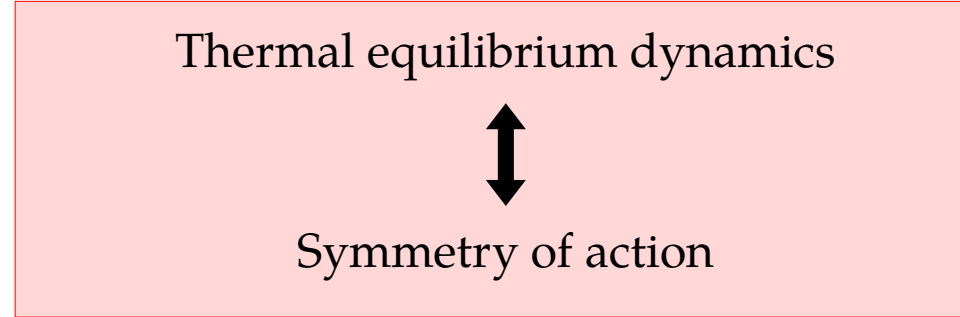


Non-equilibrium fluctuations via symmetry breaking

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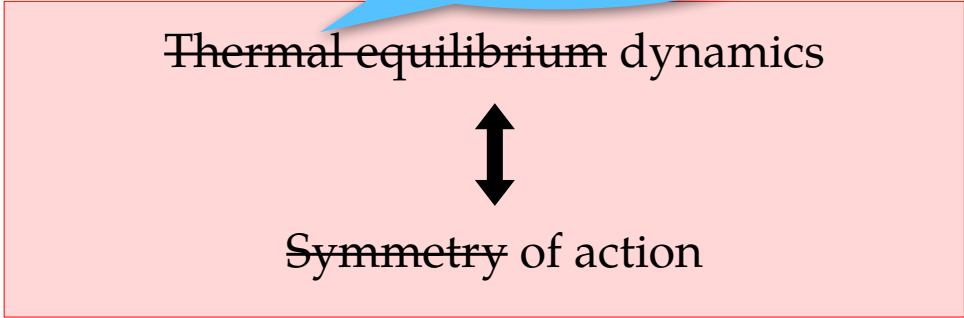


$$S[\psi] = \int dx \mathcal{L}(\psi(x), \partial_\mu \psi(x); x) \xrightarrow{\psi \mapsto \mathcal{T}_\beta[\psi]} S[\psi]$$

Ward-Takahashi identities = KMS x time reversal

$$\langle \mathcal{O}[\psi(x)] \dots \mathcal{O}[\psi(x')] \rangle = \langle \mathcal{O}[\mathcal{T}_\beta \psi(x)] \dots \mathcal{O}[\mathcal{T}_\beta \psi(x')] \rangle$$

Non-equilibrium



Symmetry breaking

$$S[\psi] = \int dx \mathcal{L}(\psi(x), \partial_u \psi(x); x) \longrightarrow S[\psi] + \Delta S$$

Ward-Takahashi identities = K

Fluctuation theorems

$$\langle \mathcal{O}[\psi(x)] \dots \mathcal{O}[\psi(x')] \rangle = \langle \mathcal{O}[\mathcal{T}_\beta \psi(x)] \dots \mathcal{O}[\mathcal{T}_\beta \psi(x')] \rangle$$

Entropy production

Thermal equilibrium dynamics



Symmetry of action

$\hbar = 0$

- Stochastic Field Theory
 - SUSY
 - MSR

Janssen (1976)
Parisi Sourlas (1982)
Aron Biroli Cugliandolo (2010)
...

Crossley Gloriosio Liu (2015)
Haehl Loganayagam Rangamani (2015)
Jensen Pinzani-Fokeeva Yarom (2017)
...

- Quantum hydrodynamics (EFT)

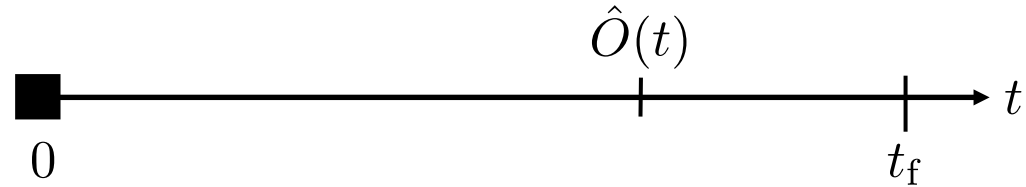
Sieberer Chiochetta Gambassi Tauber Diehl (2015)

$\hbar = 1$

- Quantum microscopies

Aron Biroli Cugliandolo (2018)
Yeo (2019)

Finite-time experiment

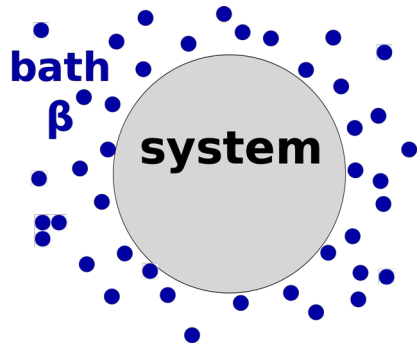


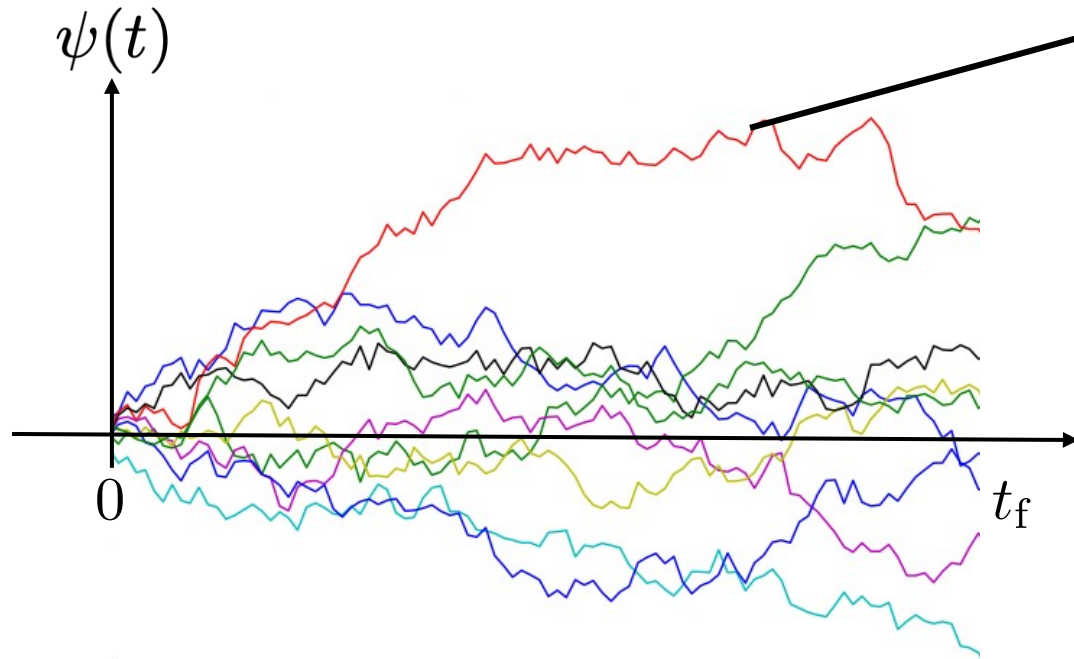
Ex: non-relativistic real scalar in 0+1 dim

$$H(t) = \frac{\pi^2}{2m} + V(\psi, \lambda(t))$$

Equilibrium conditions

- prepare Gibbs-Boltzmann statistics $\rho_0 = e^{-\beta H} / \mathcal{Z}_0$
 - isolated: wait (thermalization)
 - or couple thermostat bath
- evolve with
 - same $H = \text{const.}$
 - same bath, β





$$m\partial_t^2\psi = -V'(\psi, \lambda(t)) - \eta\partial_t\psi + \xi(t)$$

single trajectory

- work $\mathcal{W}[\psi]$
- heat $\mathcal{Q}[\psi]$
- entropy $\mathcal{S}^{\text{irr}}[\psi]$

- First law

$$\Delta\mathcal{E} = \mathcal{W}[\psi] + \mathcal{Q}[\psi]$$

- Second law

$$\Delta\mathcal{S} = \beta\mathcal{Q}[\psi] + \underbrace{\mathcal{S}^{\text{irr}}[\psi]}$$

on average $\langle\mathcal{S}^{\text{irr}}[\psi]\rangle \geq 0$

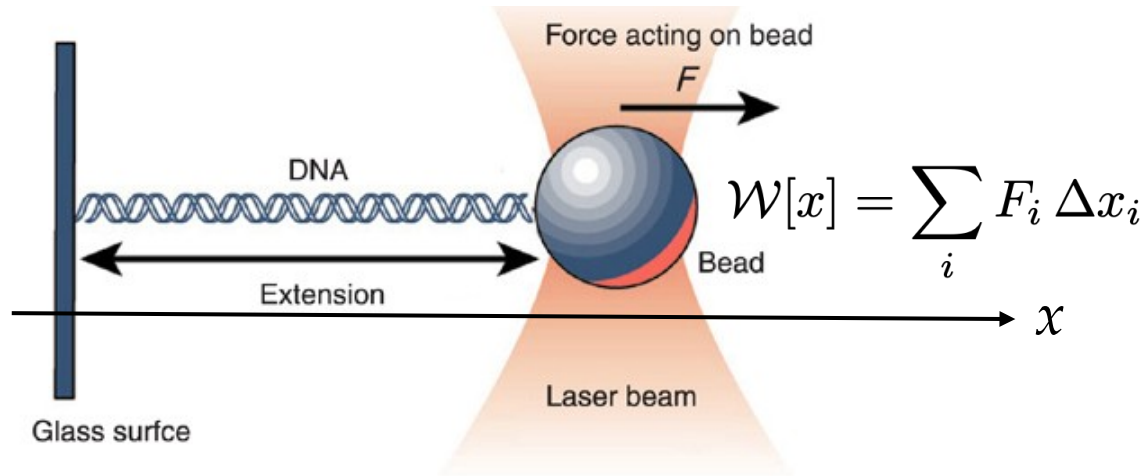
- Isolated system: Clausius inequality

$$\mathcal{W}^{\text{irr}} := \langle\mathcal{W}[\psi]\rangle - \Delta\mathcal{F} \geq 0$$

nature

Verification of the Crooks fluctuation theorem and recovery of RNA folding free energies

D. Collin^{1*}, F. Ritort^{2*}, C. Jarzynski³, S. B. Smith⁴, I. Tinoco Jr⁵ & C. Bustamante^{4,6}



Fluctuation Theorems

$$P(\mathcal{S}^{\text{irr}}) = P_r(-\mathcal{S}^{\text{irr}}) e^{\mathcal{S}^{\text{irr}}}$$

Gallavoti Cohen (1995)

Work fluctuation theorems

$$P(\mathcal{W}) = P_r(-\mathcal{W}) e^{\beta(\mathcal{W} - \Delta\mathcal{F})}$$

Crooks (1999)

$$e^{-\beta\Delta\mathcal{F}} = \langle e^{-\beta\mathcal{W}[\psi]} \rangle$$

Jarzynski (1997)

exp. review: Ciliberto (2017)

How about quantum fluctuations?

2-point projective measurement

$$\mathcal{W}_{ij} = E_i(t_f) - E_j(0)$$

Quantum jump trajectories

- work $\mathcal{W}[\psi]$
- heat $Q[\psi]$
- entropy $\mathcal{S}^{\text{irr}}[\psi]$

Mukamel, Esposito

Work fluctuation theorem

$$P(\mathcal{W}) = P_r(-\mathcal{W}) e^{\beta(\mathcal{W} - \Delta\mathcal{F})}$$

Kurchan (2000)

Stochastic equilibrium dynamics

Single particle $\psi(t)$

Model A

Model B

$$m\partial_t^2\psi = -V'(\psi) + \eta\int_0^t \xi(t')\partial_t\psi(t') + \xi(t) + \partial_t\psi = \nabla \left[\delta\mathcal{F}[\psi] + \xi(2\eta T)\delta(\mathbb{T}-t') \right]$$

Martin-Siggia-Rose-Janssen-de Dominicis

$$Z[J, \hat{J}] = \int \mathcal{D}[\psi, \hat{\psi}] e^{-S[\psi, \hat{\psi}] + \int dt \hat{J}\psi + J i \hat{\psi}}$$

$$S = S_{\text{closed}} + S_{\text{dissip}}$$

$$S_{\text{closed}} = \int dt i\hat{\psi} [m\partial_t\psi + V'(\psi)] + \beta \left(\frac{1}{2} m(\partial_t\psi(0))^2 + V(\psi(0)) \right) - \beta\mathcal{F}_0$$

$$S_{\text{dissip}} = \eta \int dt i\hat{\psi} [\partial_t\psi - T i \hat{\psi}]$$

- log(initial probability)

Equilibrium symmetry

Time-reversed fields: $\psi_r(t) := \psi(t_f - t)$

Field transformation

$$T_\beta : \begin{cases} \psi(t) & \mapsto \psi_r(t) \\ i\hat{\psi}(t) & \mapsto i\hat{\psi}_r(t) + \beta\partial_t\psi_r(t) \end{cases}$$

Invariance of the action

$$S[\psi] \xrightarrow{T_\beta} S[\psi]$$

Follow the sources

$$\int dt \hat{J}\psi + J i\hat{\psi} \xrightarrow{T_\beta} \int dt \hat{J}_r\psi + J_r [i\hat{\psi} - \beta\partial_t\psi]$$

$$Z[J, \hat{J}] \stackrel{T_\beta}{=} Z[J_r, \hat{J}_r + \beta\partial_t J_r]$$

Chatelain (2003)

Aron Biroli Cugliandolo (2010)

Aron Barci Cugliandolo G-Arenas Lozano (2014)

Aron ... Lozano (2016)

→ Fluctuation-dissipation theorems

$$R(t, t') = R(t_f - t, t_f - t') + \beta\partial_{t'} C(t_f - t, t_f - t')$$

$$C(t, t') = \langle \psi(t)\psi(t') \rangle$$

$$R(t, t') = \langle \psi(t)i\hat{\psi}(t') \rangle$$

Symmetry breaking out of equilibrium

Time-dependent non-equilibrium drive: $V(\psi) \rightarrow V(\psi, \lambda)$

Variation of the action under \mathcal{T}_β

$$S_{\text{closed}}[\psi, \lambda] \xrightarrow{\mathcal{T}_\beta} S_{\text{closed}}[\psi, \lambda_r] + \beta \int dt \dot{\lambda}_r \partial_\lambda V(\psi, \lambda_r) - \beta \Delta \mathcal{F}_r \quad \lambda_r(t) := \lambda(t_f - t)$$

$$S_{\text{dissip}}[\psi] \xrightarrow{\mathcal{T}_\beta} S_{\text{dissip}}[\psi]$$

Measure of irreversibility

Broken Ward identities

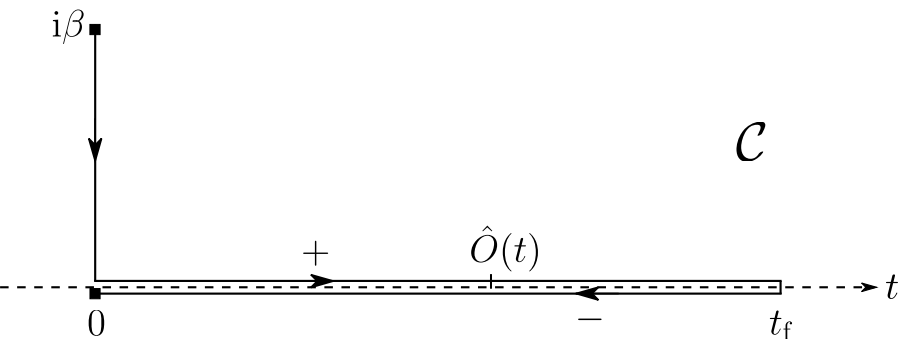
$$\langle O(\psi(t)) \rangle = \langle O(\psi_r(t)) e^{-\beta(\mathcal{W}_r[\psi] - \Delta \mathcal{F}_r)} \rangle_r$$

→ **Work fluctuation theorems**

$$\frac{P(\mathcal{W})}{P_r(-\mathcal{W})} = e^{\beta(\mathcal{W}[\psi] - \Delta \mathcal{F})}$$

Quantum equilibrium dynamics

$$\text{ex: } H = \frac{\pi^2}{2m} + V(\psi) \quad [H, \Theta] = 0$$



Kadanoff-Baym contour

$$Z[J^+, J^-] = \int \mathcal{D}[\psi^+, \psi^-] e^{iS[\psi] + \int dt J^+ \psi^+ - J^- \psi^-}$$

$$\times \mathcal{Z}_0^{-1} \langle \psi^+(0) | e^{-\beta H} | \psi^-(0) \rangle \langle \psi^-(t_f) | \psi^+(t_f) \rangle$$

$$S[\psi] = \int dt \mathcal{L}[\psi^+(t)] - \int dt \mathcal{L}[\psi^-(t)]$$

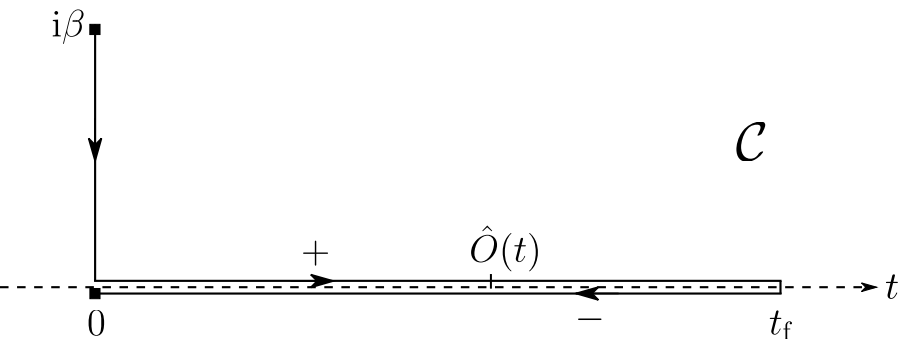
Tentative symmetry transformation

$$\psi^\pm(t) \mapsto \psi^\pm(t_f - t \pm i\beta/2)$$

manifestly not symmetric



Complex-time reparametrization



change of frame: R_θ

$$|\psi(t)\rangle \mapsto e^{i\theta(t)H(t)} |\psi(t)\rangle \quad \theta(t) \in \mathbb{C}$$

$$\mathbf{I} = \int d\psi(t) e^{+i\theta(t)H(t)} |\psi(t)\rangle \langle \psi(t)| e^{-i\theta(t)H(t)}$$

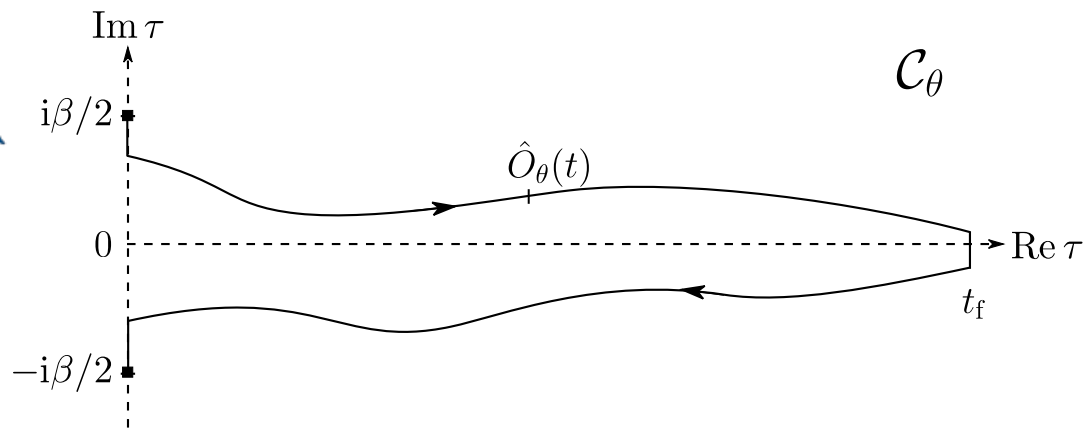
$$H(t) \mapsto e^{-i\theta(t)H(t)} [H(t) - i\partial_t] e^{+i\theta(t)H(t)}$$

new close-time contour

$$t \mapsto \tau = t + \theta(t) \quad \psi(t) \mapsto \psi(\tau)$$

Follow the sources

$$\text{ex: } \int dt J^+ \psi^+ \xrightarrow{\mathcal{R}_\theta} \int dt J^+ e^{-i\theta H} \psi^+ e^{+i\theta H}$$

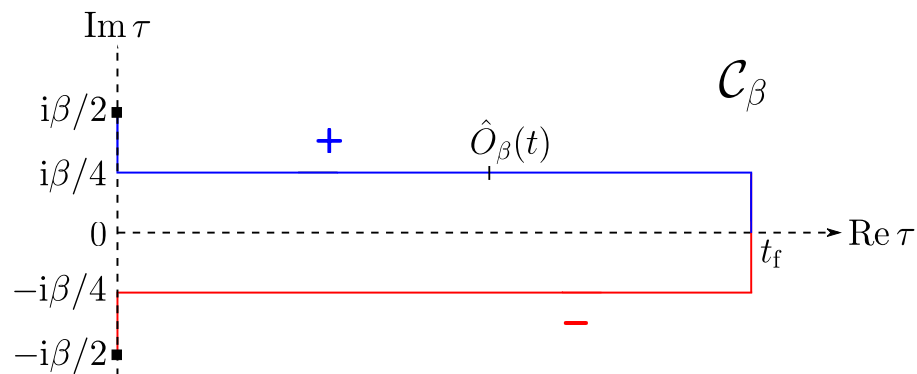


Symmetry of equilibrium

Equilibrium conditions: $\hat{H} = \text{const}$, $\hat{\rho}_0 = e^{\beta(\mathcal{F} - \hat{H})}$

Choice of frame/contour: $\theta(t) = \pm i\beta/4$

Action: $iS[\psi] = i \int_{\mathcal{C}_\beta} d\tau \mathcal{L}[\psi(\tau)] + \beta\mathcal{F}_0$



Field transformation

$$T_\beta : \psi^\pm(\tau) \mapsto \psi^\pm(t_f - \tau \pm i\beta/2)$$

Invariance of the action

$$S[\psi] \xrightarrow{T_\beta} S[\psi]$$

KMS symmetry

$$Z[J^+(t), J^-(t)] \stackrel{\mathcal{R}_\beta^{-1} \circ \mathcal{T}_\beta \circ \mathcal{R}_\beta}{=} Z[J^+(t_f - t + i\beta/2), J^-(t_f - t - i\beta/2)]$$

Fluctuation-dissipation theorem

$$G^{-+}(t, t') \stackrel{\mathcal{R}_\beta^{-1} \circ \mathcal{T}_\beta \circ \mathcal{R}_\beta}{=} G^{+-}(t_f - t' + i\beta/2, t_f - t - i\beta/2)$$

Symmetry of equilibrium

Open quantum systems

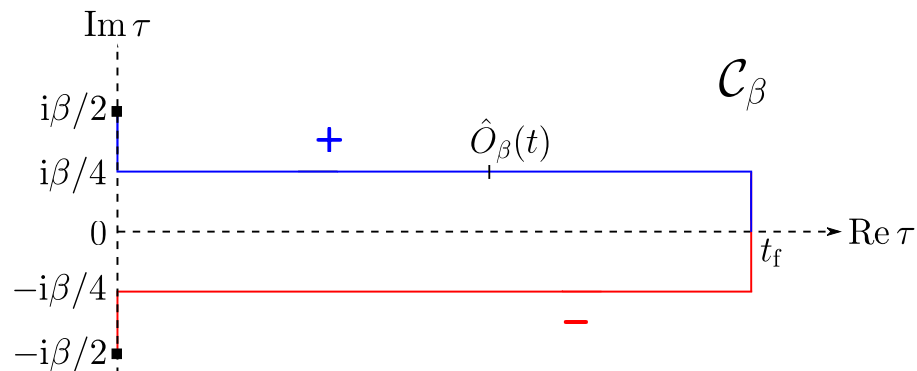
Caldeira-Leggett bath

$$H = H_S + H_{SB} + H_B$$

$$H_S = \frac{\pi^2}{2m} + V(\psi)$$

$$H_B = \sum_i \frac{p_i^2}{2m_i} + \frac{1}{2} m \omega_i^2 q_i^2$$

$$H_{SB} = \frac{1}{2} \sum_i c_i (q_i - \psi)^2$$



$$S_{\text{dissip}} = \int_{C_\beta} d\tau \int^\tau d\tau' \psi^+(\tau) K(\tau - \tau') \psi^+(\tau') \\ + \int d\tau \int^\tau d\tau' \psi^-(\tau) K(\tau - \tau') \psi^-(\tau') \\ + \int d\tau \int d\tau' \psi^+(\tau) K(\tau - \tau') \psi^-(\tau')$$

$$K(\tau) := \int d\epsilon J(\epsilon) [\coth(\beta\epsilon/2) \cosh(i\epsilon\tau) - \sinh(i\epsilon\tau)]$$

Field transformation

$$T_\beta : \psi^\pm(\tau) \mapsto \psi^\pm(t_f - \tau \pm i\beta/2)$$

Invariance of the action

$$S_{\text{dissip}}[\psi] \xrightarrow{T_\beta} S_{\text{dissip}}[\psi]$$

Symmetry breaking out of equilibrium

Non-equilibrium time-dependent drive: $H \mapsto H(t)$

ex: $H(t) = \frac{\pi^2}{2m} + V(\psi, \lambda(t)) \quad H(t) = \int dx h(x) + \lambda(t)O_0$

R_β Thermal frame: non-Hermitian evolution

$$H \mapsto \begin{cases} H^+ = H - e^{\beta H/4} i\partial_t e^{-\beta H/4} \\ H^- = (H^+)^\dagger \end{cases}$$

$$H \mapsto H_r = \frac{\pi^2}{2m} + V(\psi, \lambda_r) \quad \lambda_r(t) := \lambda(t_f - t)$$

R_β^{-1} \dots

Variation of the Kadanoff-Baym action under $\mathcal{R}_\beta^{-1} \circ \mathcal{T}_\beta \circ \mathcal{R}_\beta$

$$iS[\psi] \mapsto iS_r[\psi] + \beta(\Delta\mathcal{F}_r - \Omega_r[\psi])$$

Measure of irreversibility

with $\beta \Omega_r[\psi] := \int dt \langle \psi^+(t) | \dot{\sigma} | \psi^+(t) \rangle + \int dt \langle \psi^-(t) | \dot{\sigma} | \psi^-(t) \rangle$

$\psi \in \mathbb{R}$

Symmetry breaking out of equilibrium

$$\dot{\sigma} := -e^{\beta H/2} \left[\frac{d}{dt} e^{-\beta H} \right] e^{\beta H/2}$$

- Semiclassical expansion

$$\dot{\sigma} \simeq \beta \partial_t H + \frac{1}{3!} (\beta/2)^2 [H, H, [\beta \partial_t H]] + \dots$$

- Spatially extended

$$\xi \lesssim v_{\text{LR}} \times \beta \hbar$$

- Non-locality in time

$$\langle \dot{\sigma} \rangle = \int_{-\beta \hbar/2}^{\beta \hbar/2} d\tau \text{Tr}[\partial_t H \rho(t + i\tau)] \rightarrow \tau \sim \beta \hbar$$

- Weak Clausius-like inequality

$$\int dt \langle \dot{\sigma} \rangle \geq \beta \int dt \langle \partial_t H \rangle \geq \beta \Delta \mathcal{F}$$



Symmetry breaking out of equilibrium

Non-equilibrium time-dependent drive: $H \mapsto H(t)$

ex: $H(t) = \frac{\pi^2}{2m} + V(\psi, \lambda(t))$ $H(t) = \int dx h(x) + \lambda(t)O_0$

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$$H \mapsto H_r = \frac{\pi^2}{2m} + V(\psi, \lambda_r)$$

$$\lambda_r(t) := \lambda(t_f - t)$$

• • •

Work fluctuation theorem

$$P(\mathcal{W}) = P_r(-\mathcal{W}) e^{\beta(\mathcal{W} - \Delta\mathcal{F})}$$

with $\mathcal{W}_{ij} = E_i(t_f) - E_j(0)$

finite-time fluctuations: care for boundaries

	classical thermal	quantum
Equilibrium symmetry	<p>MSR</p> $S[\psi] \xrightarrow{\mathcal{T}_\beta} S[\psi]$	<p>Kadanoff-Baym</p> $\mathcal{R}_\beta \downarrow \uparrow \mathcal{R}_\beta^{-1}$ <p>'Thermal' frame $S[\psi] \xrightarrow{\mathcal{T}_\beta} S[\psi]$</p>
	<div style="border: 1px solid black; padding: 5px; display: inline-block;">Fluctuation-Dissipation Theorems</div>	
Non-equilibrium symmetry breaking	$S[\psi] \mapsto S_r[\psi] - \beta(\Delta\mathcal{F}_r - \mathcal{W}_r[\psi])$ <p style="text-align: center;">work Fluctuation Theorems</p>	$iS[\psi] \mapsto iS_r[\psi] + \beta(\Delta\mathcal{F}_r - \Omega_r[\psi])$ <p style="text-align: center;">non-Hermitian dynamics work Fluctuation Theorems</p>